

Primordial non-Gaussianity & the Galaxy Bispectrum

Cosmological non-Gaussianity Workshop
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w/ Martin Crocce & Vincent Desjacques

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Matter correlators

a bit of Perturbation Theory ...

$$P = P_0 + P_G^{loop}[P_0] + P_{NG}^{loop}[P_0, B_0]$$

matter power spectrum

Linear power spectrum

Gravity-induced contributions
(depending on P_0 alone)

Additional gravity-induced contributions
present *only* for NG initial conditions (B_0)

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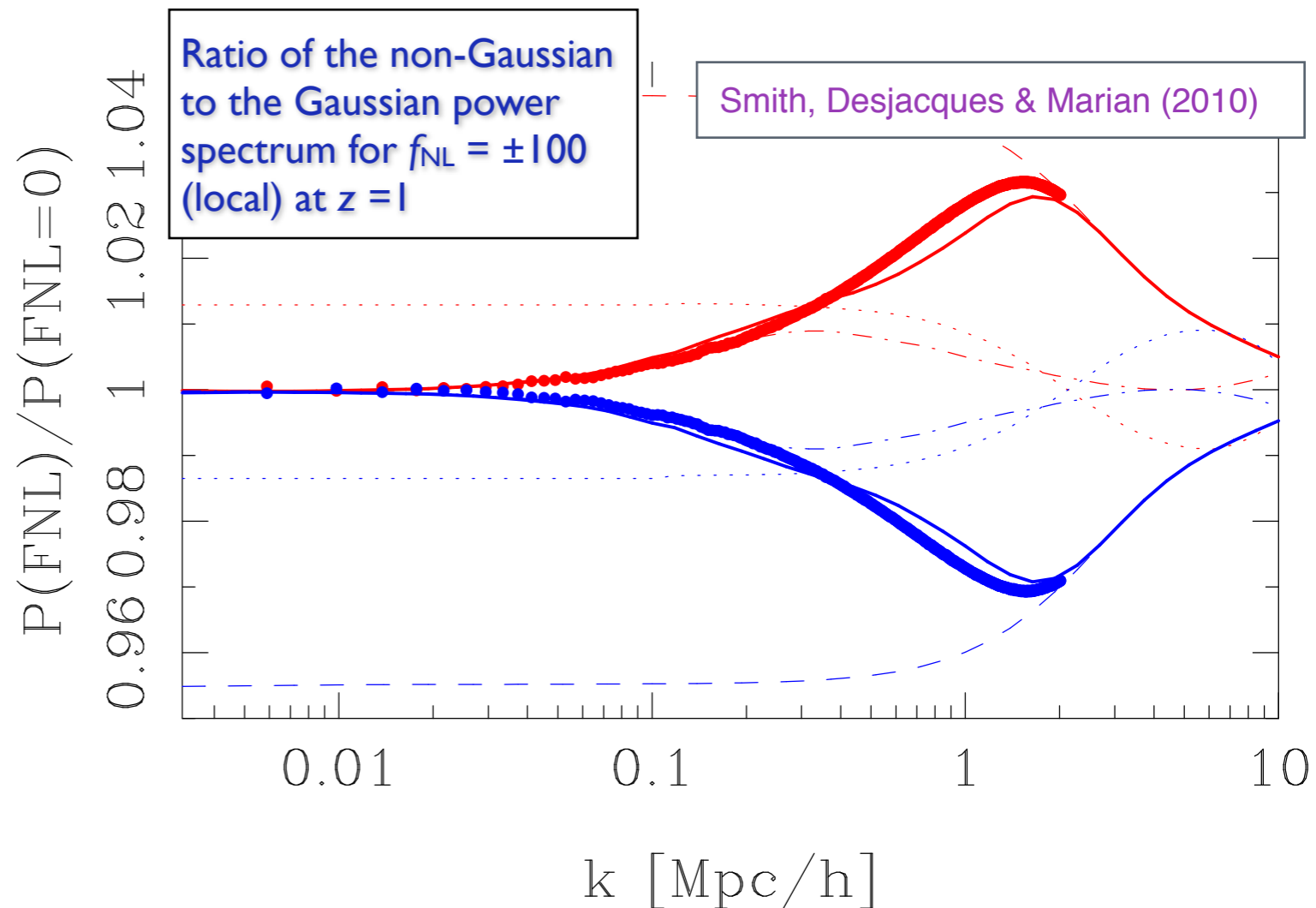
matter power spectrum

Linear power spectrum

Gravity-induced contributions (depending on P_0 alone)

Additional gravity-induced contributions present *only* for NG initial conditions (B_0)

Few percent effect at small scales for allowed values of f_{NL}



Matter correlators

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& bispectrum

Primordial component

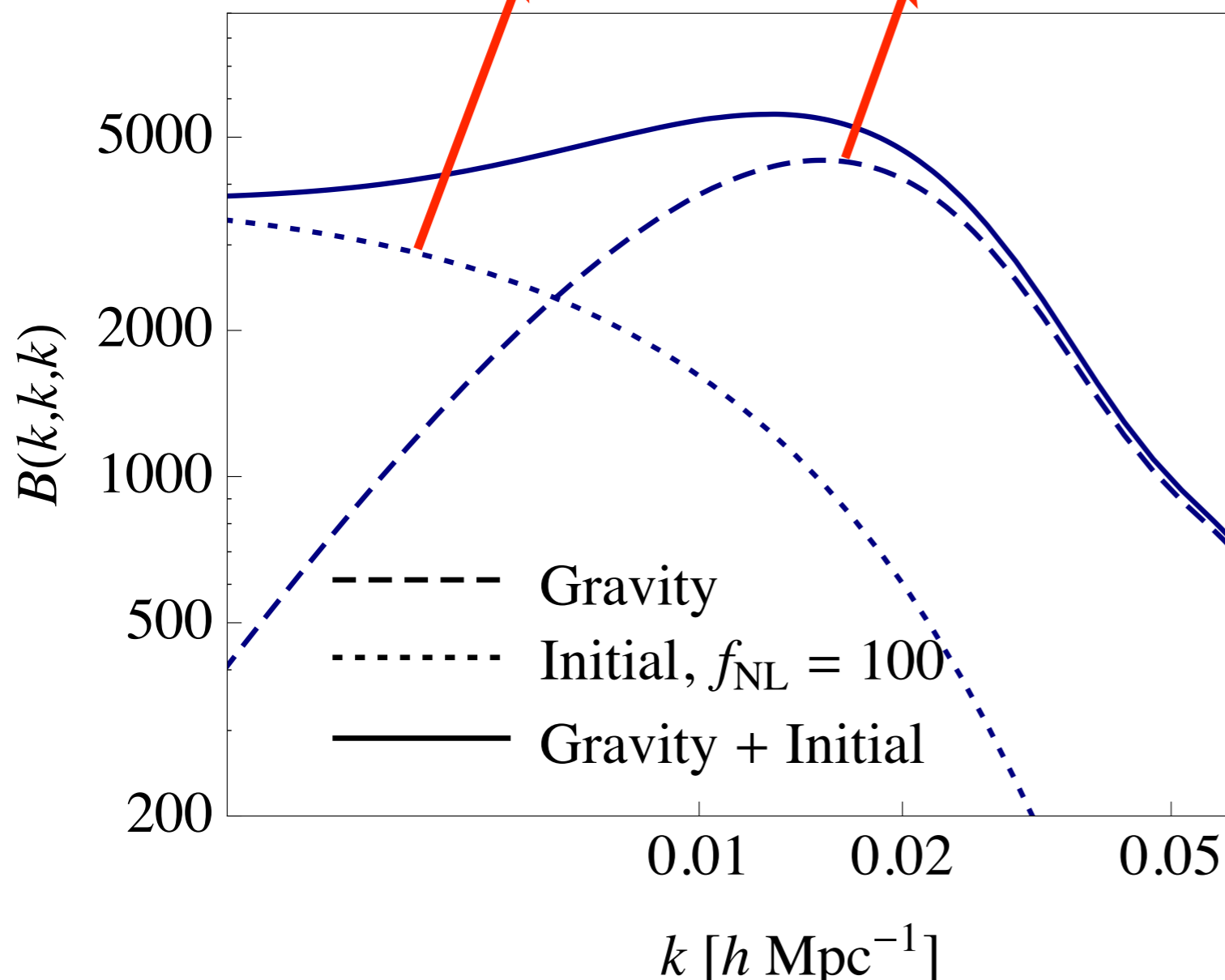
The matter bispectrum and PNG: *large scales*

At large scales

$$B(k_1, k_2, k_3) \simeq B_0 + B_G^{tree}[P_0]$$

↓
Primordial
component

↓
Gravity-induced
component



Equilateral configurations of the matter bispectrum

$$\frac{B_0(k, k, k)}{B_G^{tree}(k, k, k)} \underset{k \rightarrow 0}{\sim} \frac{f_{NL}}{D(z)k^2}$$

The primordial component has a different dependence on scale than the gravity-induced one

This is true for almost all models (local, equilateral, orthogonal ...)

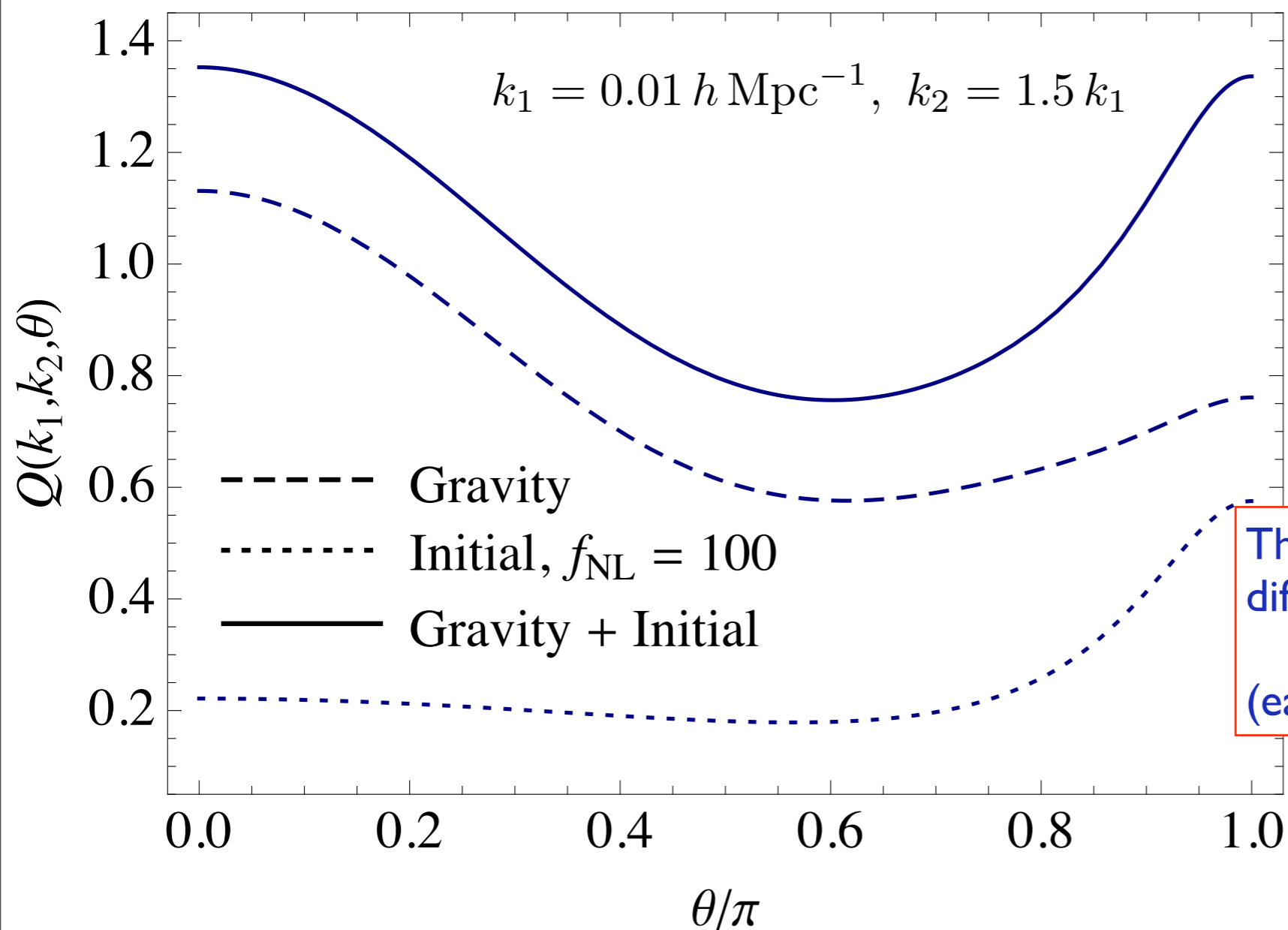
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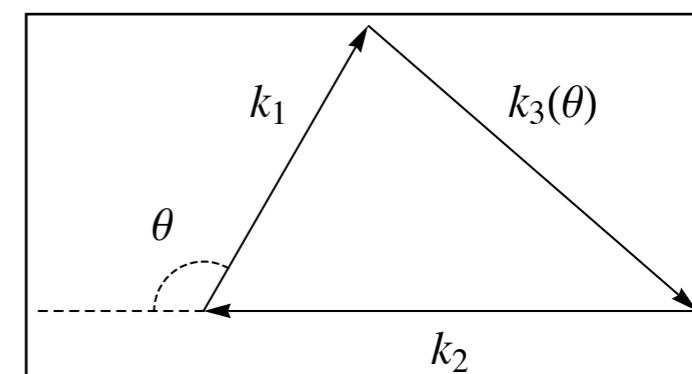
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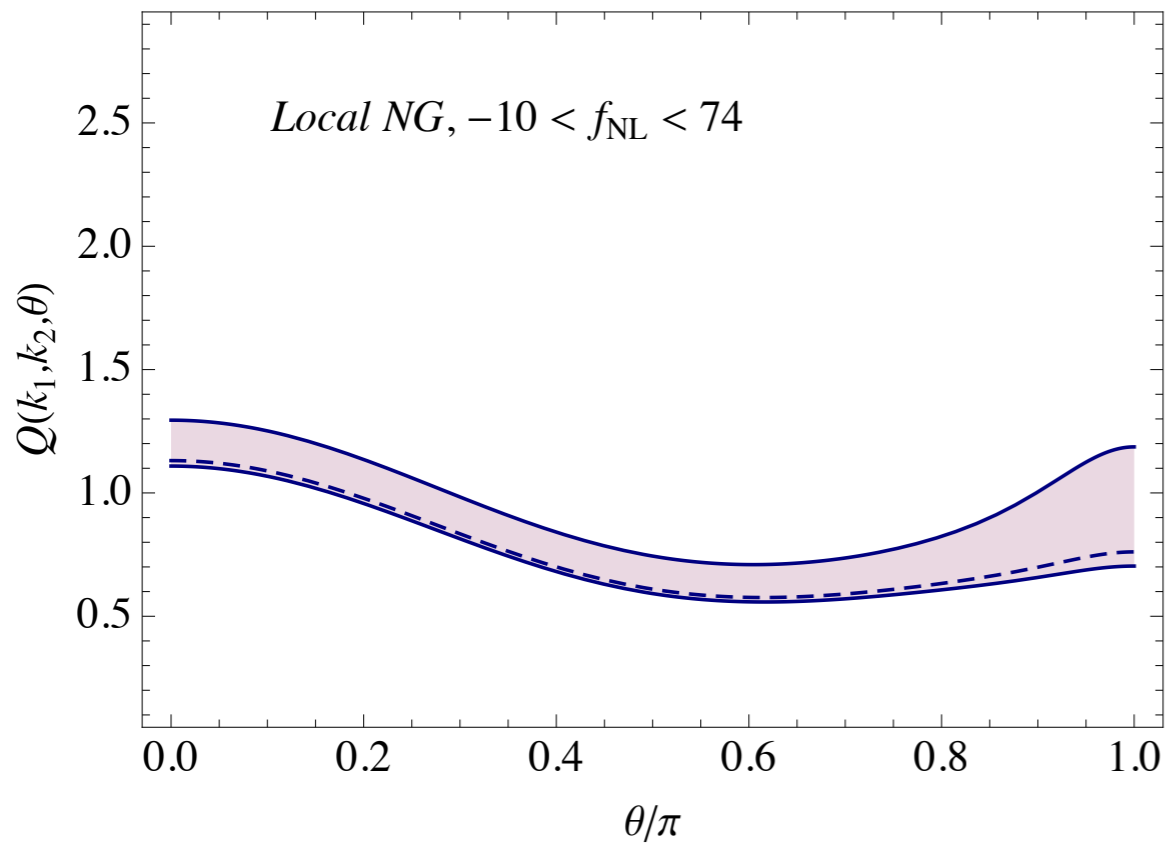


$$Q = \frac{B}{P(k_1)P(k_2) + \text{cyc.}}$$

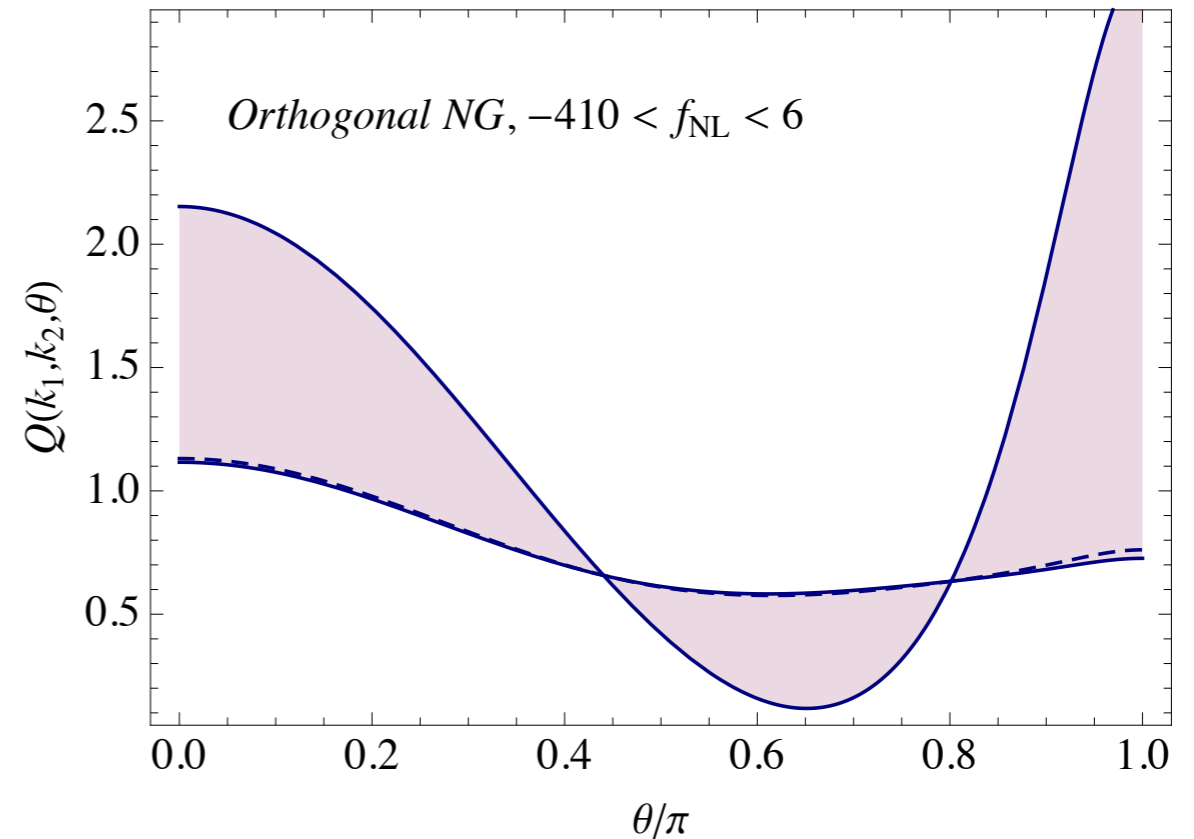
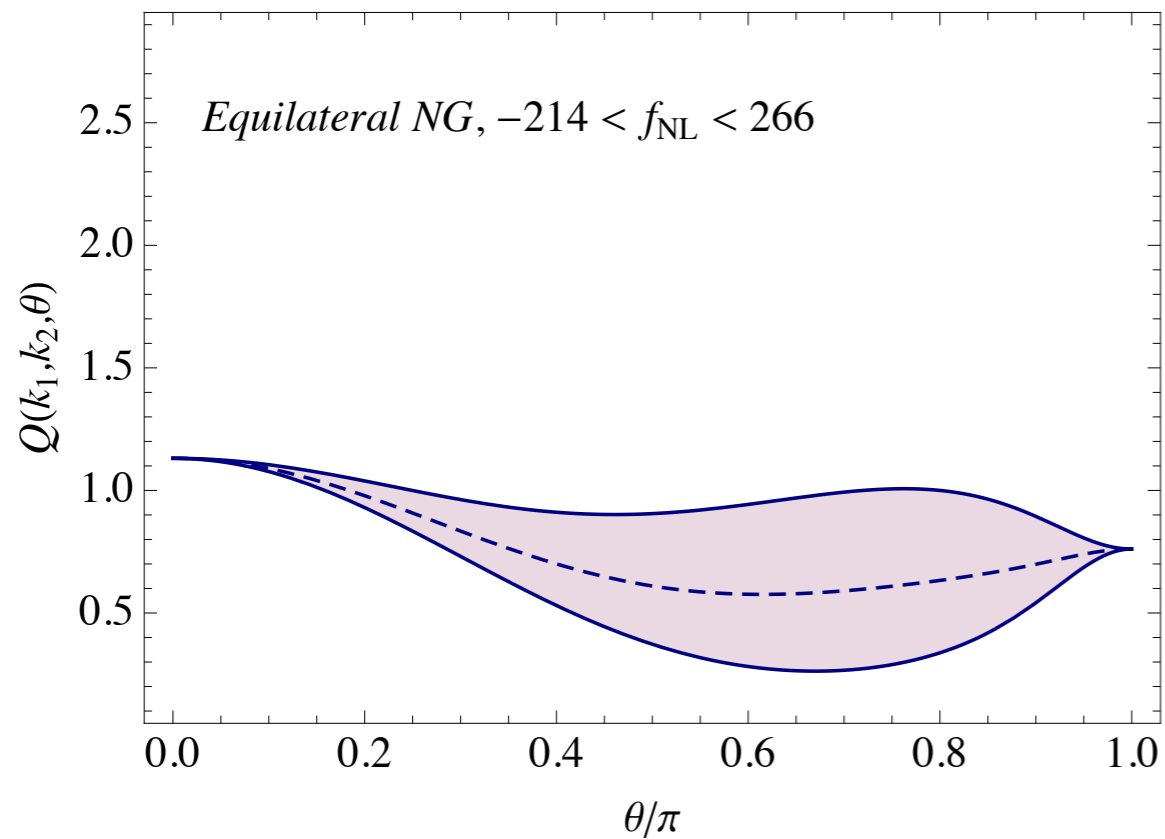


The primordial component has a different shape dependence
 (each model has its own, of course)

The matter bispectrum and PNG: *large scales*

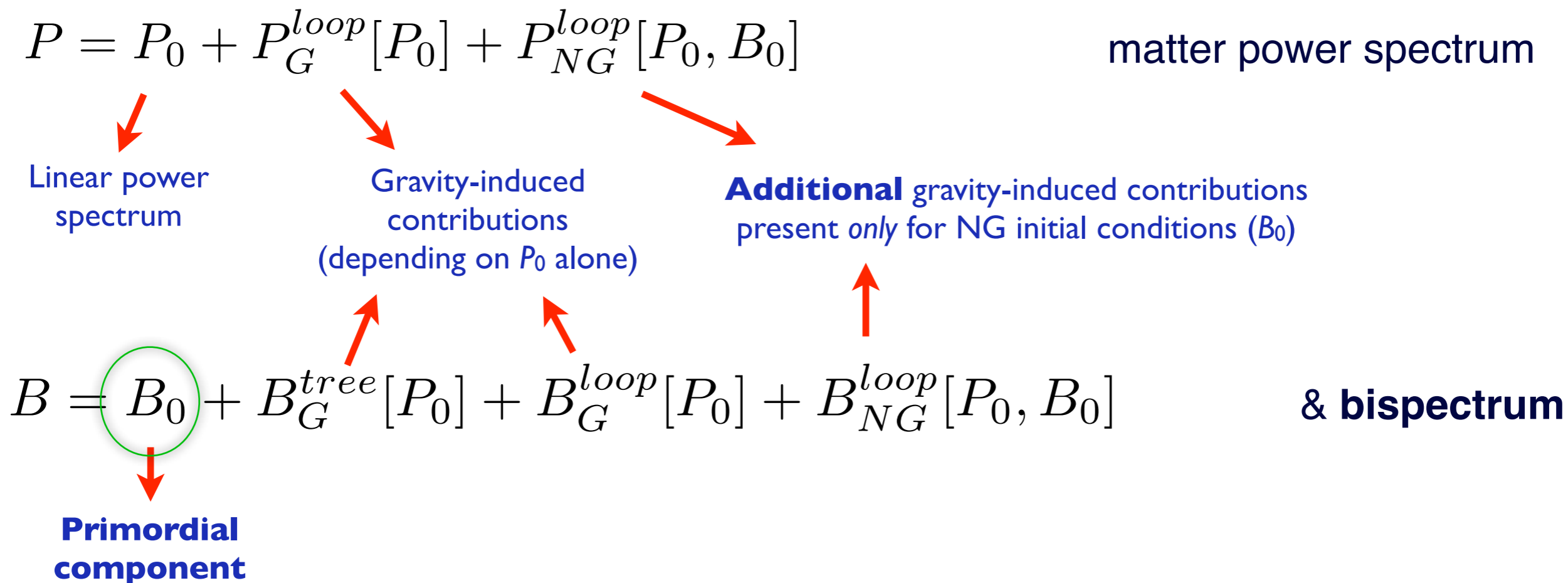


Current CMB constraints for different models of non-Gaussianity as uncertainties on generic configurations of the matter bispectrum, $B \simeq B_0 + B_G^{\text{tree}}[P_0]$



Matter correlators

a bit of Perturbation Theory ...



If B_0 was the *only effect* of NG initial conditions on the LSS then future, large volume surveys ($\sim 100 \text{ Gpc}^3$) could provide:

$$\Delta f_{\text{NL}}^{\text{local}} < 5 \text{ and } \Delta f_{\text{NL}}^{\text{eq}} < 10$$

ES & Komatsu (2007)

Matter correlators

a bit of Perturbation Theory ...

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matter power spectrum

Linear power spectrum

Gravity-induced contributions (depending on P_0 alone)

Additional gravity-induced contributions present *only* for NG initial conditions (B_0)

$$B = B_0 + B_G^{tree}[P_0] + B_G^{loop}[P_0] + B_{NG}^{loop}[P_0, B_0]$$

& bispectrum

Primordial component

Nonlinear corrections are also affected by the initial conditions!

There is a **significant effect** of NG initial conditions of about 5-15% on all triangles, at **small scales** and at **late times** for $f_{NL} = 100$

The matter bispectrum and PNG: *small scales*

$$B = B_0 + B_G^{tree}[P_0] + B_G^{loop}[P_0] + B_{NG}^{loop}[P_0, B_0]$$

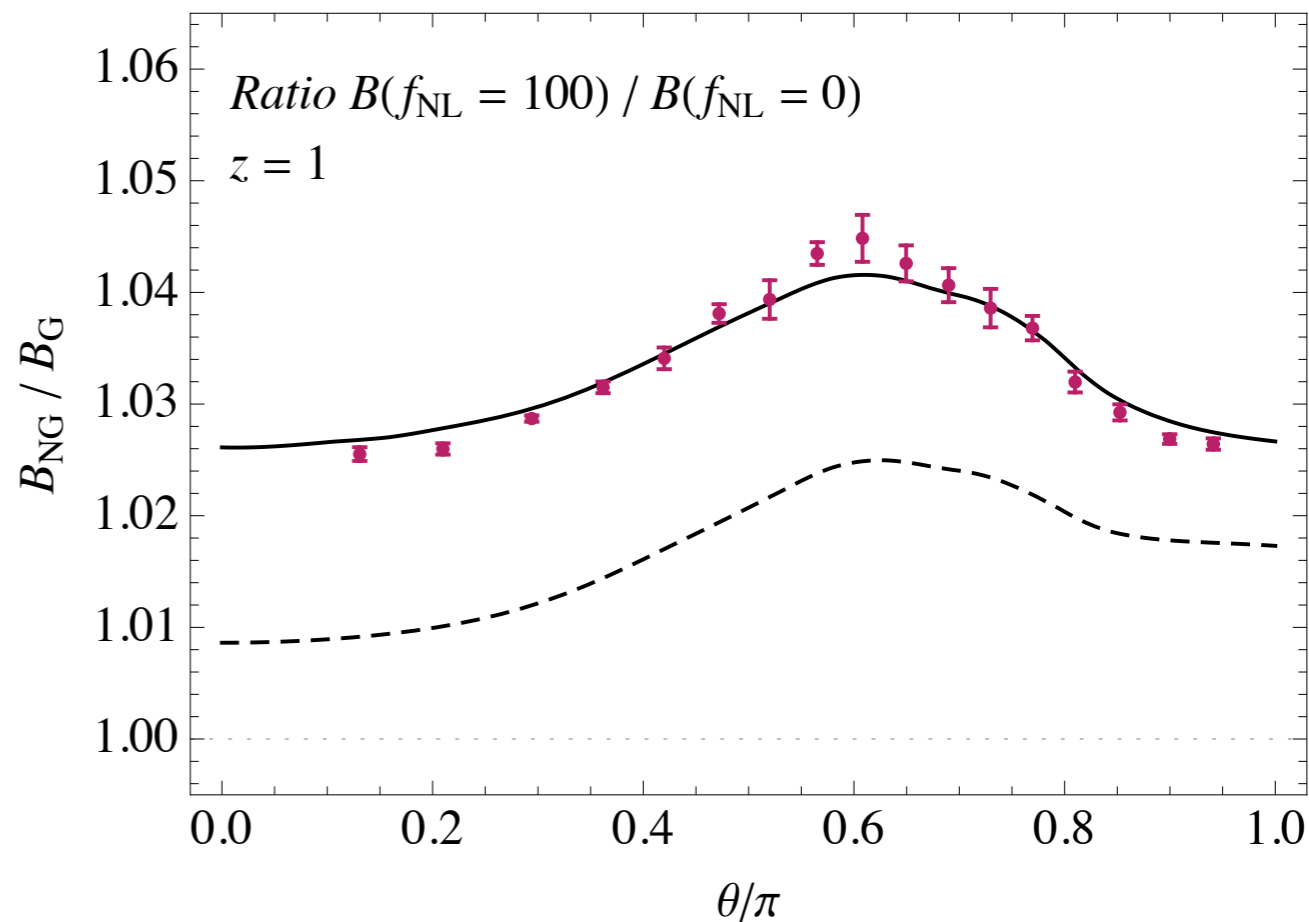
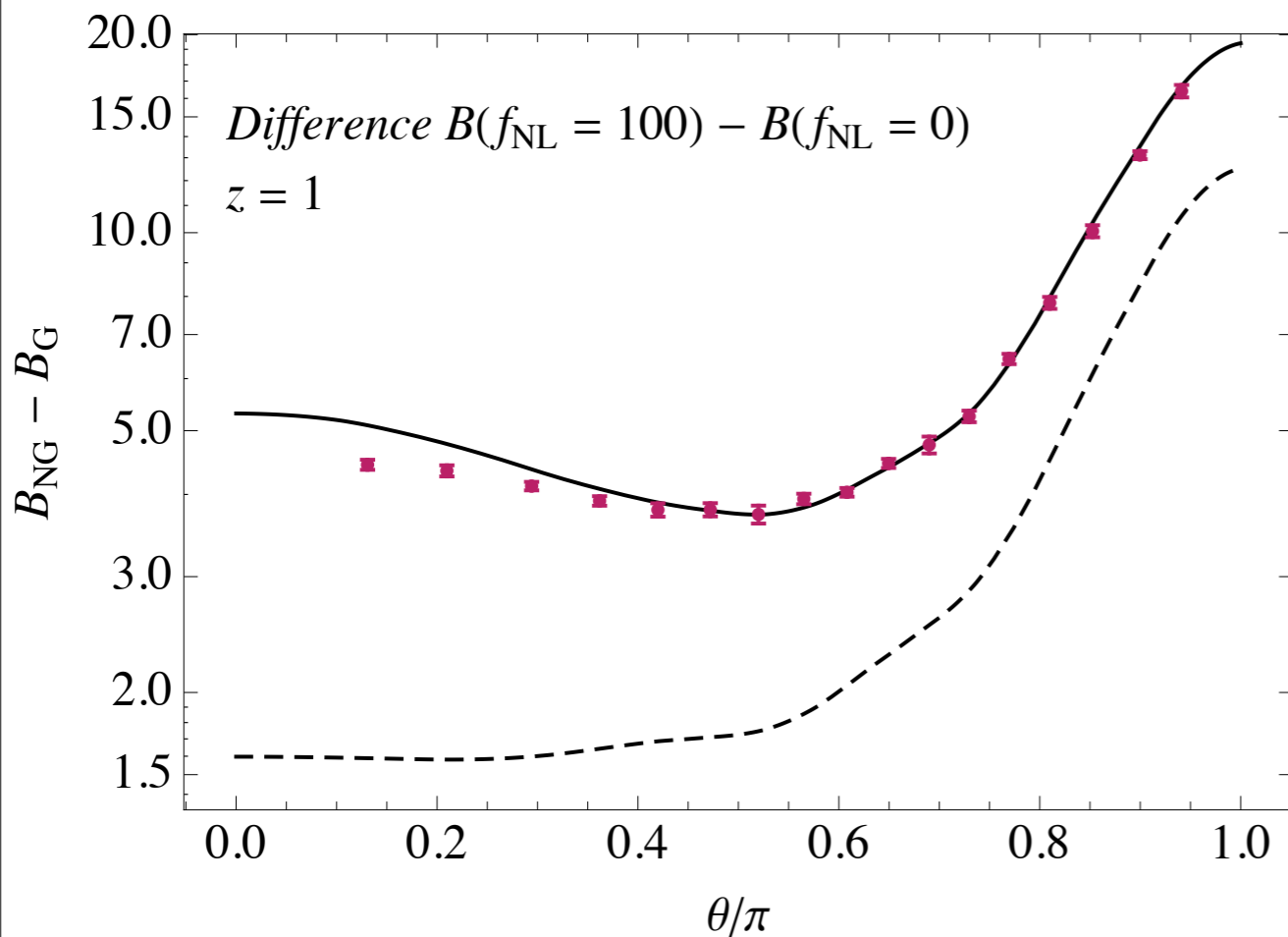
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Primordial
component

↓ Gravity-induced
contributions

↓ **Additional** gravity-induced contributions
present for NG initial conditions (B_0)

Generic configurations $B(k_1, k_2, \theta)$
as a function of θ
with $k_1 = 0.1 h/\text{Mpc}$, $k_2 = 1.5 k_1$

ES (2009)
ES, Crocce & Desjacques (2010)



The matter bispectrum and PNG: *small scales*

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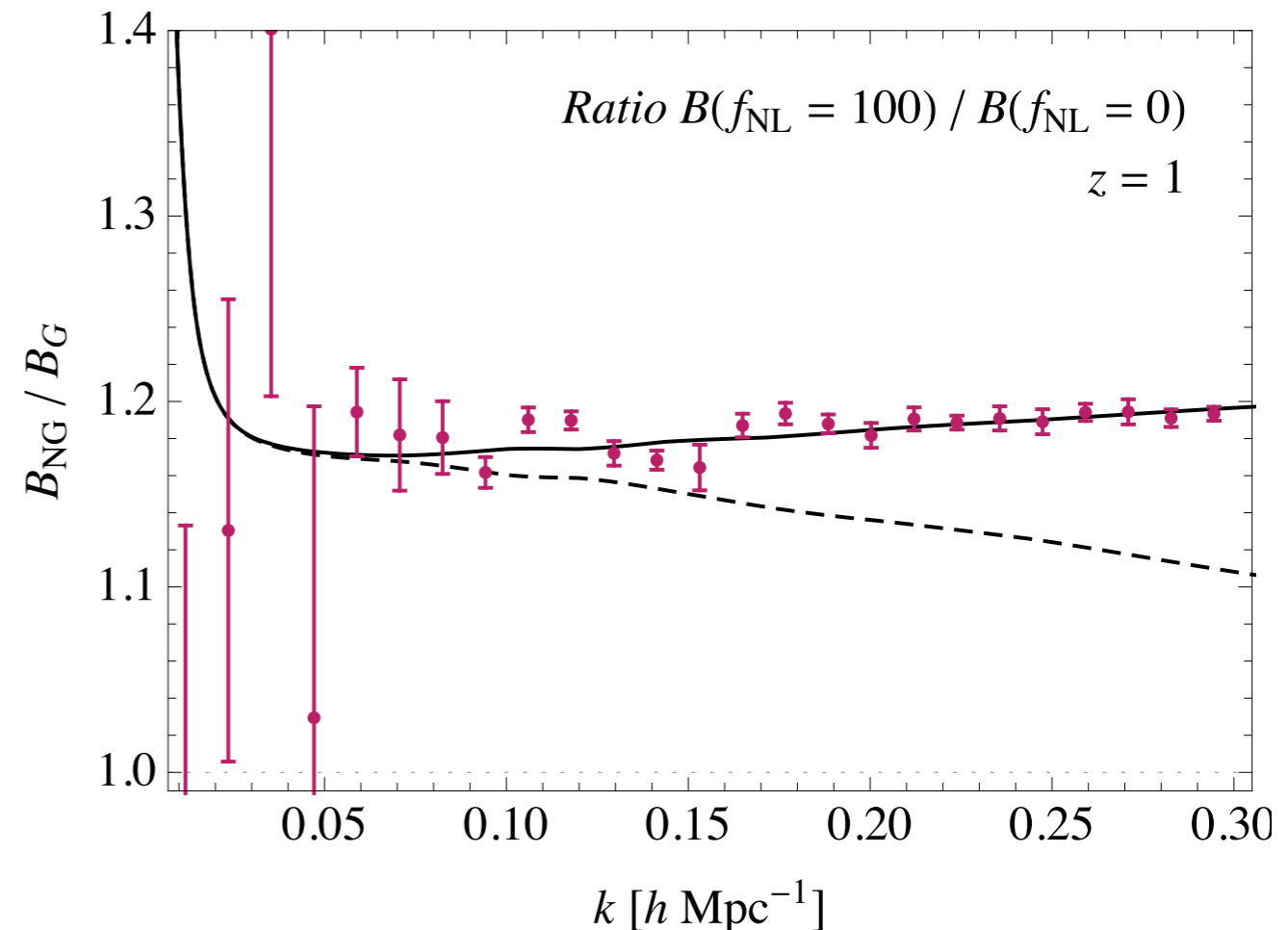
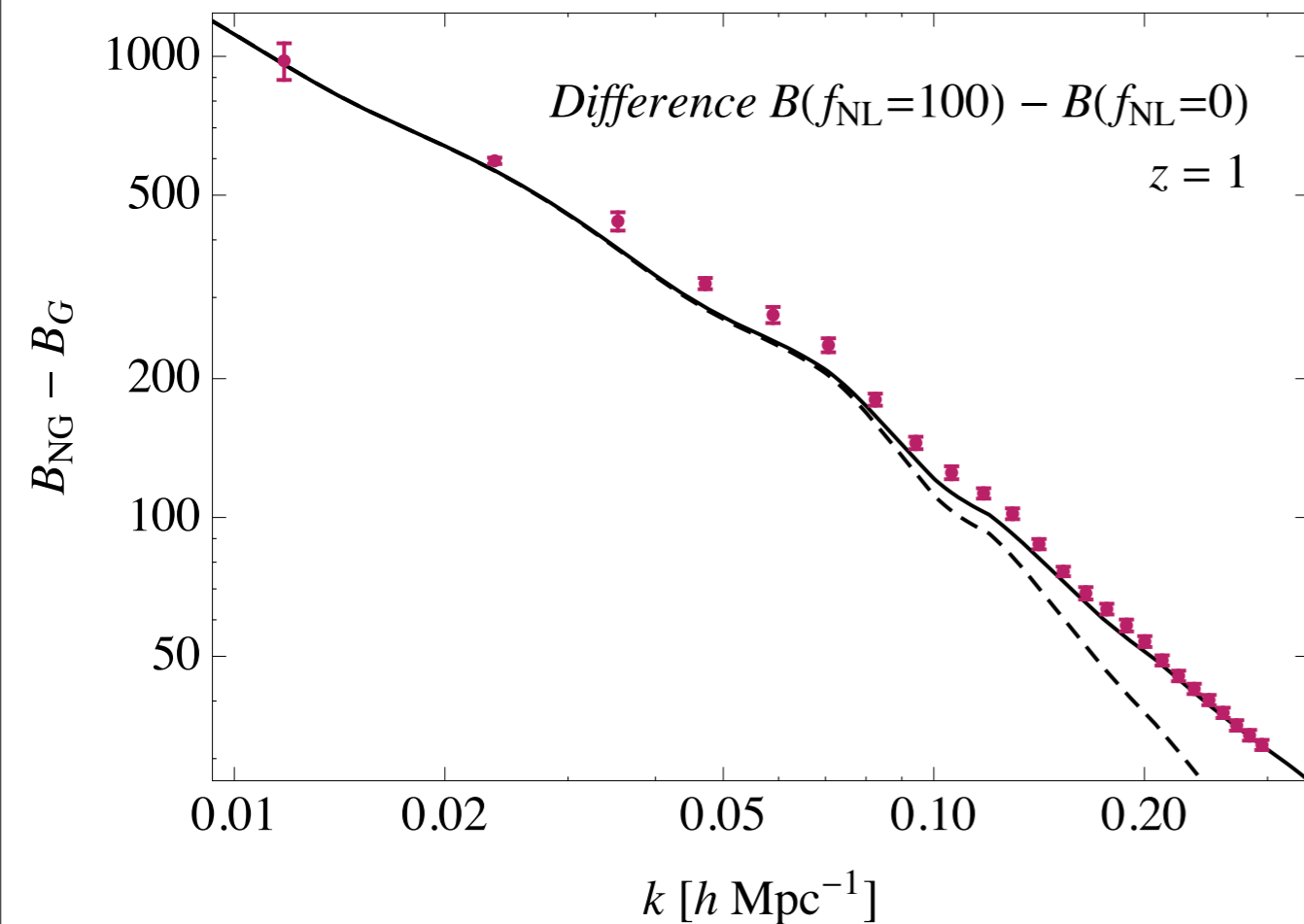
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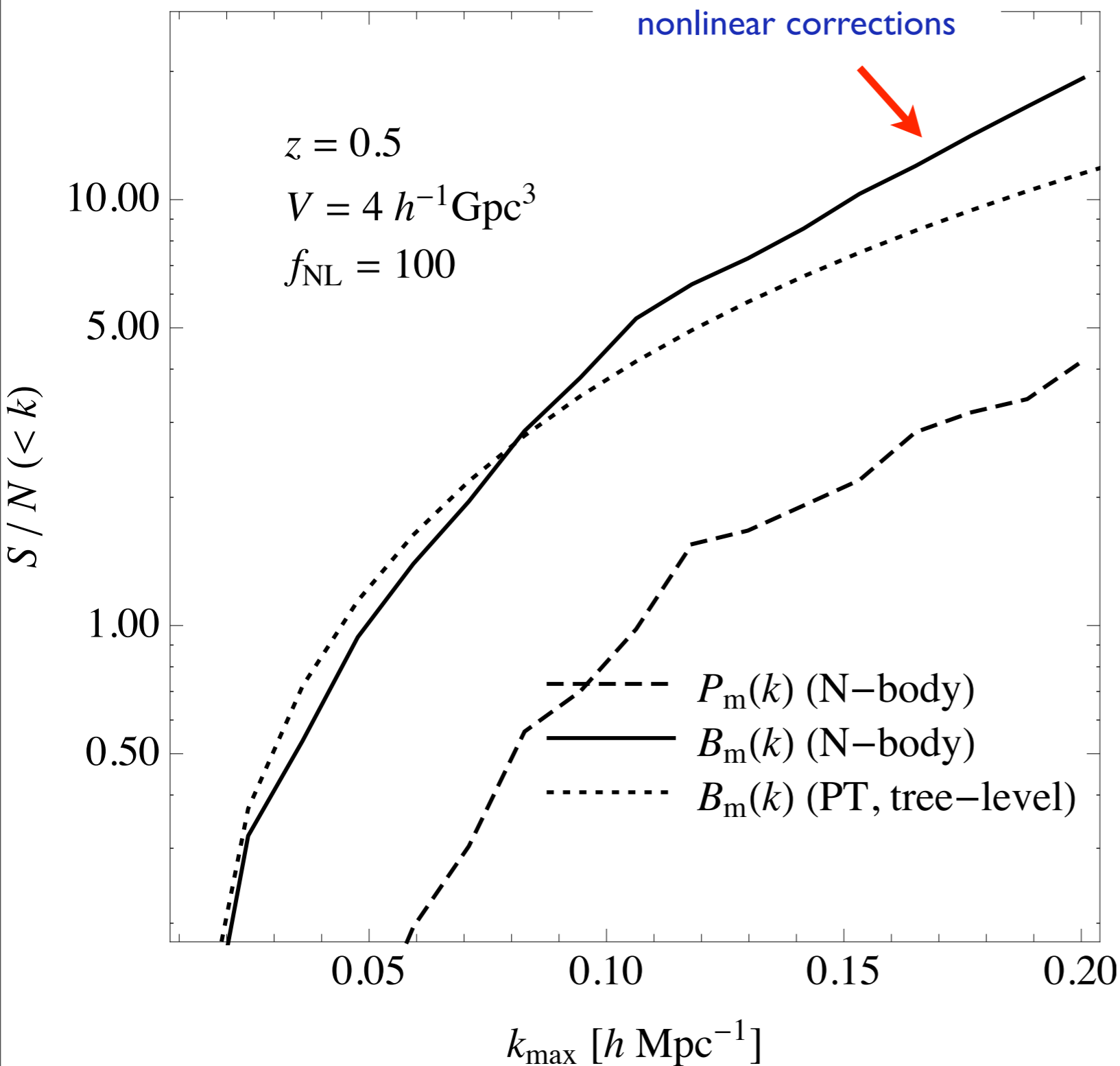
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Squeezed configurations $B(\Delta k, k, k)$
as a function of k with $\Delta k = 0.01 h/\text{Mpc}$

ES (2009)
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Matter Power Spectrum vs Matter Bispectrum



Cumulative signal-to-noise for the effect of NG initial conditions.

Sum of all configurations up to k_{max}

$$\left(\frac{S}{N}\right)_P^2 = \sum_k^{k_{\text{max}}} \frac{(P_{\text{NG}} - P_G)^2}{\Delta P^2}$$

$$\left(\frac{S}{N}\right)_B^2 = \sum_{k_1, k_2, k_3}^{k_{\text{max}}} \frac{(B_{\text{NG}} - B_G)^2}{\Delta B^2}$$

ES, Crocce & Desjacques
(in preparation)

Effects of PNG on the galaxy power spectrum

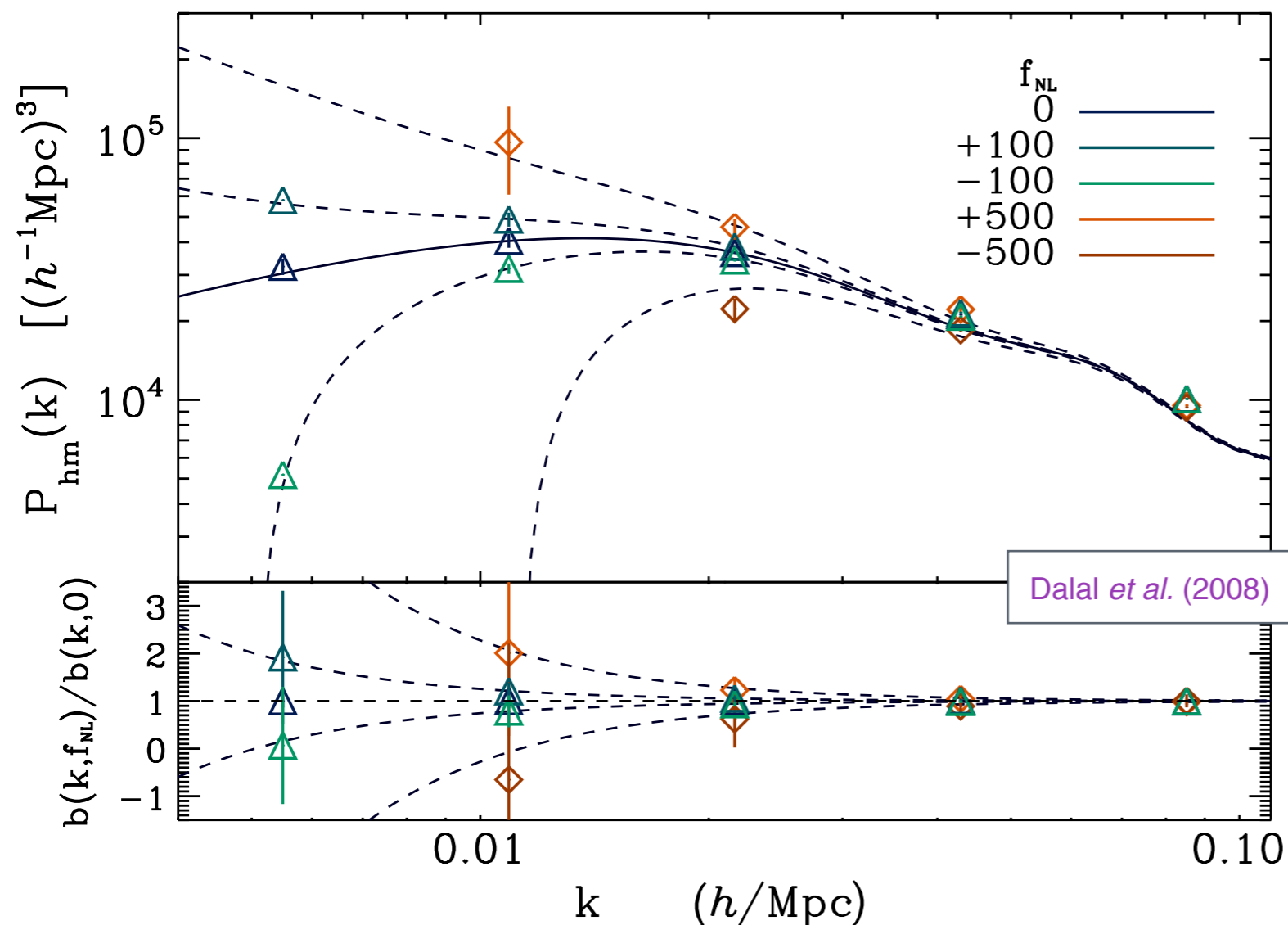
Dalal et al. (2008):

$$\delta_g(\vec{k}) = [b_1 + \Delta b_1(f_{NL}, k)] \delta_{\vec{k}} + \dots \rightarrow P_g(k) = [b_1 + \Delta b_1(f_{NL}, k)]^2 P(k)$$

↓
“Gaussian”
bias

↓
Scale-dependent correction
due to local non-Gaussianity

$$\Delta b_{1,NG}(f_{NL}, k) \sim \frac{f_{NL}}{D(z) k^2}$$



Effects of PNG on the **galaxy bispectrum**

Clearly, the effect on galaxy bias affects as well the **galaxy bispectrum**

$$B_g(k_1, k_2, k_3) = b_1^3 B(k_1, k_2, k_3) + b_1^2 b_2 P(k_1) P(k_2) + 2 \text{ perm.} + \dots$$

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Clearly, the effect on galaxy bias affects as well the **galaxy bispectrum**

$$B_g(k_1, k_2, k_3) = \overset{b_{1,G} + \Delta b_{1,NG}(f_{NL}, k)}{\uparrow} B(k_1, k_2, k_3) + \overset{b_{2,G} + \Delta b_{2,NG}(f_{NL}, \vec{k}_1, \vec{k}_2)}{\uparrow} b_1^2 P(k_1) P(k_2) + 2 \text{ perm.} + \dots$$

Scale-dependent
bias corrections

$$\Delta b_{1,NG}(f_{NL}, \vec{k}) = \Delta b_{1,si}(f_{NL}) + \Delta b_{1,sd}(f_{NL}, b_{1,G}, \vec{k})$$

$$\Delta b_{2,NG}(f_{NL}, \vec{k}_1, \vec{k}_2) = \Delta b_{2,si}(f_{NL}) + \Delta b_{2,sd}(f_{NL}, b_{1,G}, b_{2,G}, \vec{k}_1, \vec{k}_2)$$

Giannantonio & Porciani (2010)
Baldauf, Seljak & Senatore (2010)

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Scale-dependent bias corrections

$$B = B_0 + B_G^{tree}[P_0] + B_G^{loop}[P_0] + B_{NG}^{loop}[P_0, B_0]$$

Primordial component
(large scales)

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Primordial component (large scales) Effect on nonlinear evolution (small scales)

$$\Delta b_{1,NG}(f_{NL}, \vec{k}) = \Delta b_{1,si}(f_{NL}) + \Delta b_{1,sd}(f_{NL}, b_{1,G}, \vec{k})$$

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Giannantonio & Porciani (2010)
Baldauf, Seljak & Senatore (2010)

- We test this model in N-body simulations with local NG initial conditions

$$\langle \delta\delta\delta_h \rangle = \delta_D(\vec{k}_{123}) B_{mmh}$$

- We fit **all** triangular configurations up to $k = 0.07 \text{ h/Mpc}$ for **$b_{1,G}$** , **$b_{2,G}$** , **$\Delta b_{1,G}$** and **$\Delta b_{2,G}$**

$$\langle \delta_h\delta_h\delta_h \rangle = \delta_D(\vec{k}_{123}) B_h$$

$$P_h \rightarrow b_{1,G}, \Delta b_{1,si}$$

$$B_{h,G} \rightarrow b_{2,G}$$

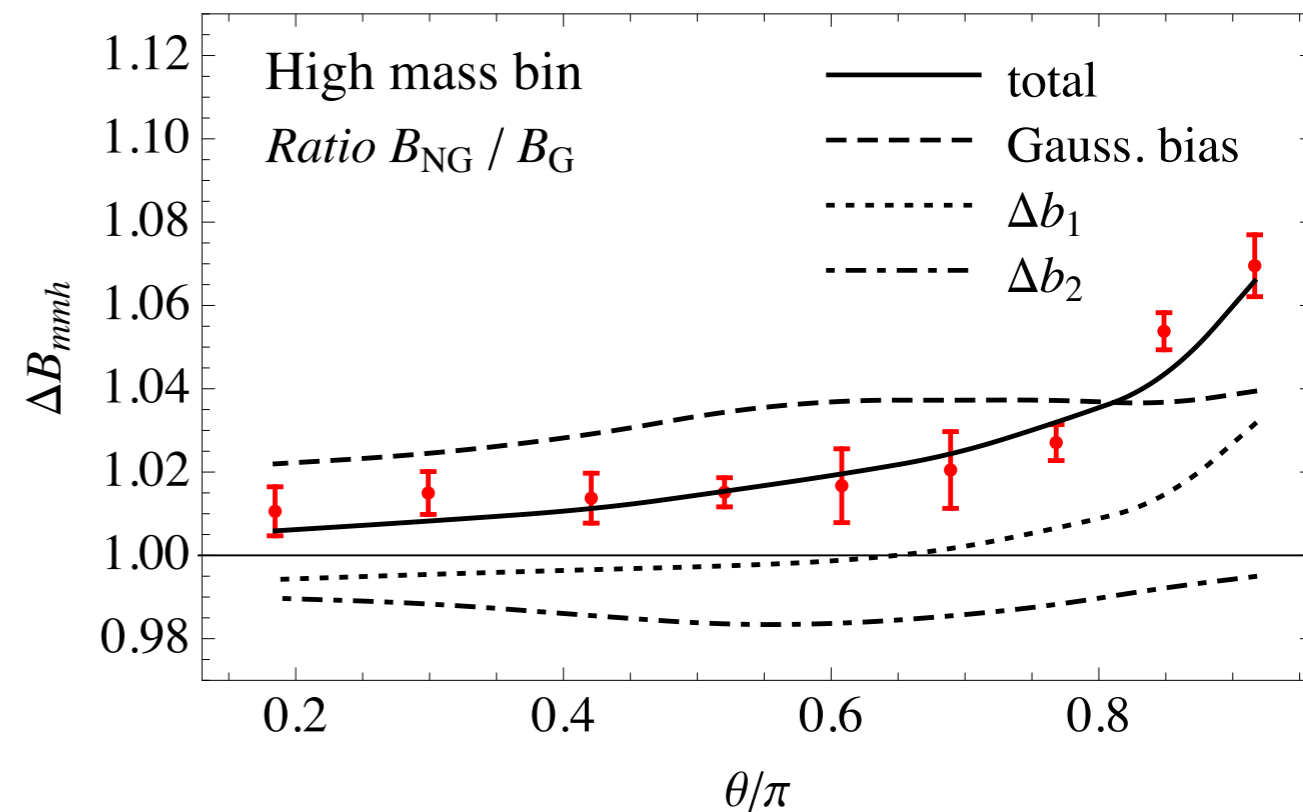
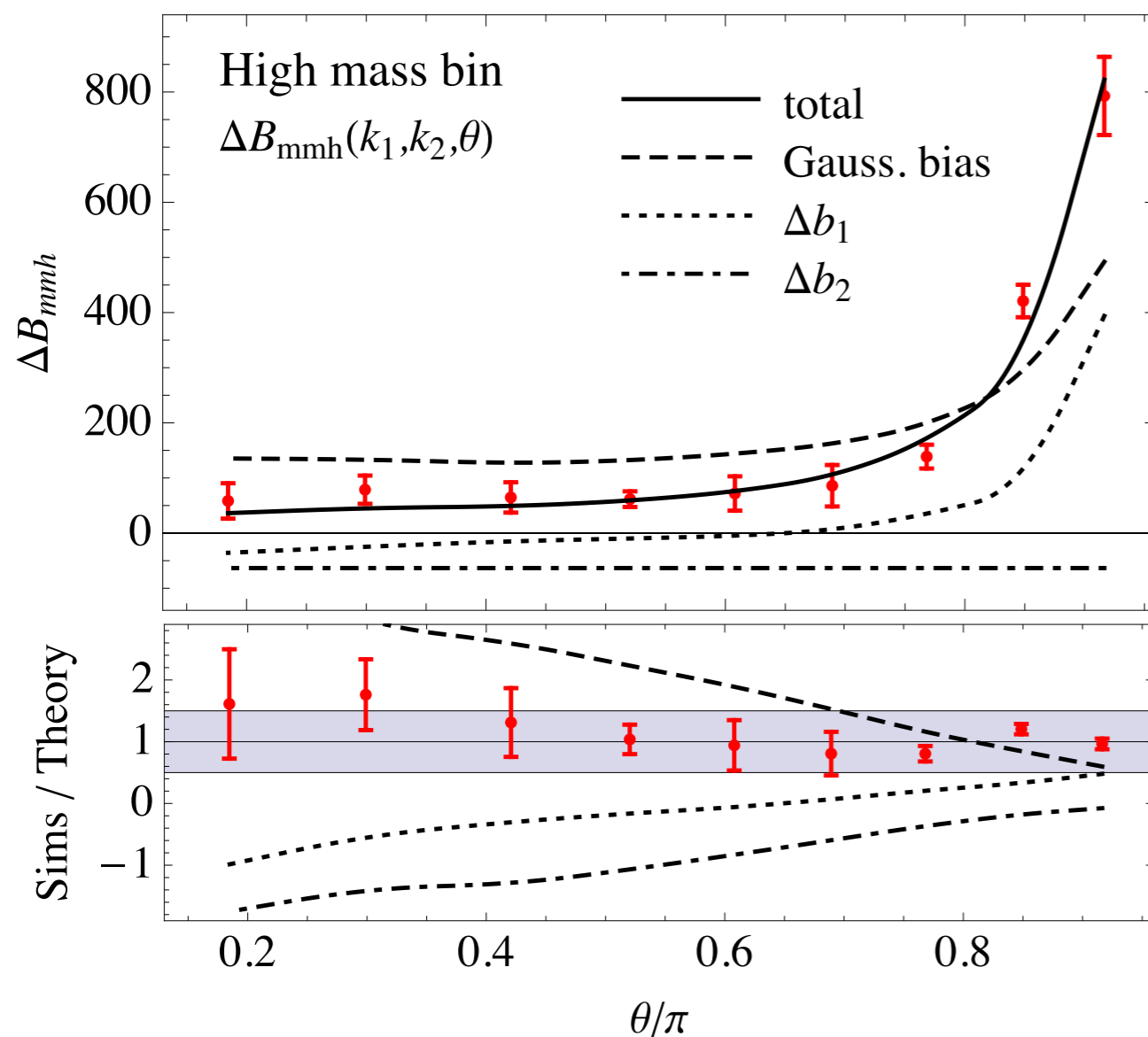
$$\Delta B_{h,NG} \rightarrow \Delta b_{2,si}$$

Effects of PNG on the galaxy bispectrum

Matter-matter-halo bispectrum:

$$B_{mmh}(k_1, k_2; k_3) = b_1(f_{NL}, k) B(k_1, k_2, k_3) + b_2(f_{NL}, k_1, k_2) P(k_1) P(k_2)$$

Generic configurations $B(k_1, k_2, \theta)$
as a function of θ
with $k_1 = 0.1 \text{ h/Mpc}$, $k_2 = 1.5 k_1$



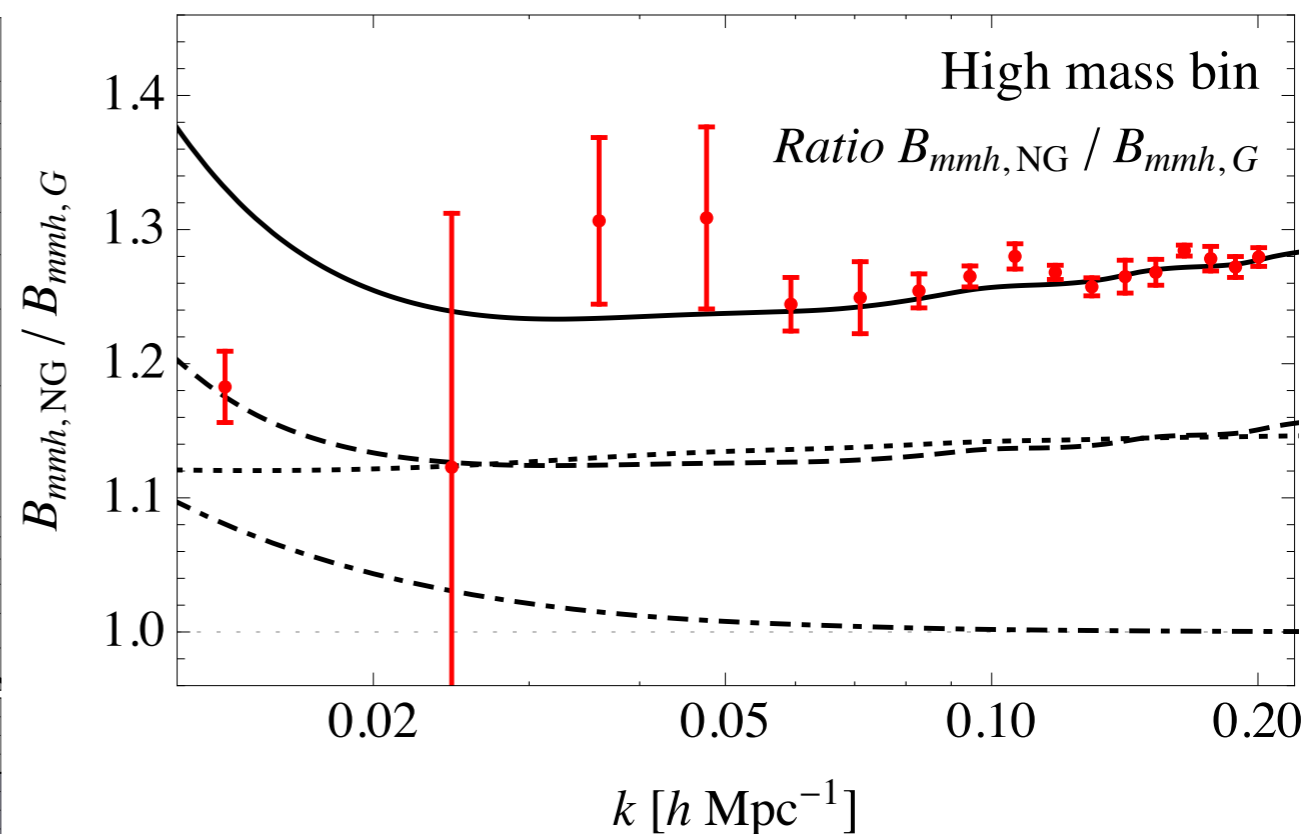
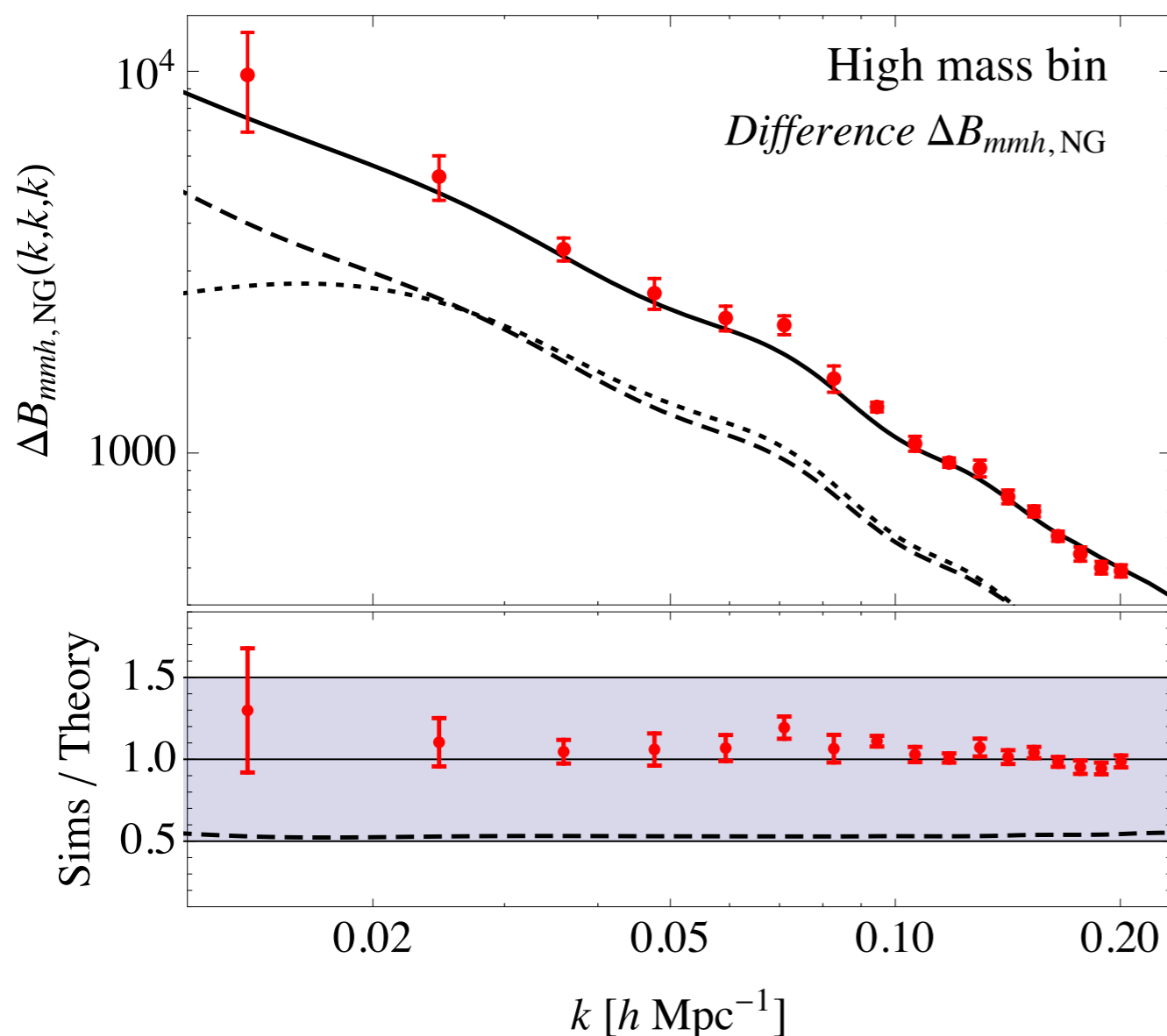
ES, Crocce & Desjacques (*in preparation*)

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Matter-matter-halo bispectrum:

$$B_{mmh}(k_1, k_2; k_3) = b_1(f_{NL}, k) B(k_1, k_2, k_3) + b_2(f_{NL}, k_1, k_2) P(k_1) P(k_2)$$

Squeezed configurations $B(\Delta k, k, k)$
as a function of k with $\Delta k = 0.01 h/\text{Mpc}$



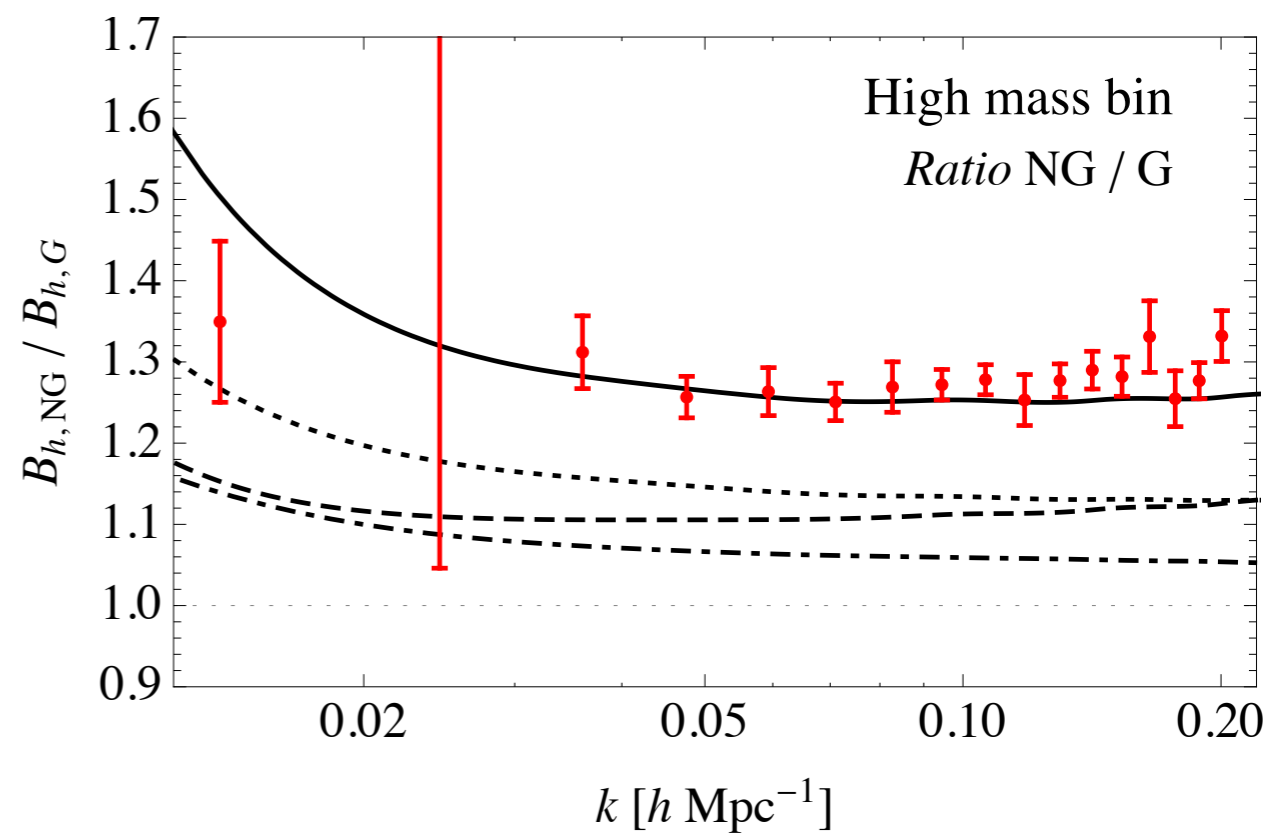
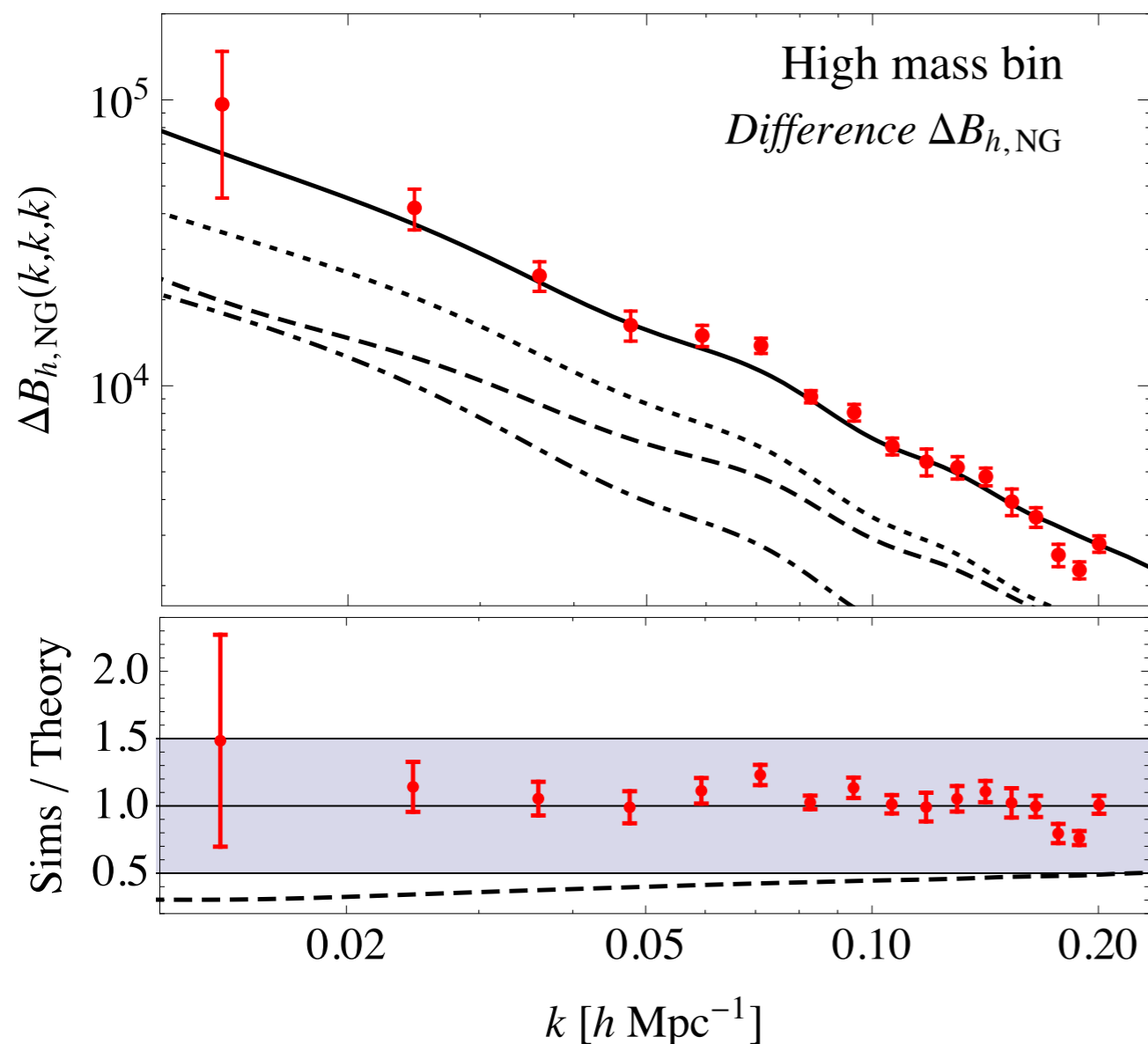
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Effects of PNG on the galaxy bispectrum

Halo bispectrum:

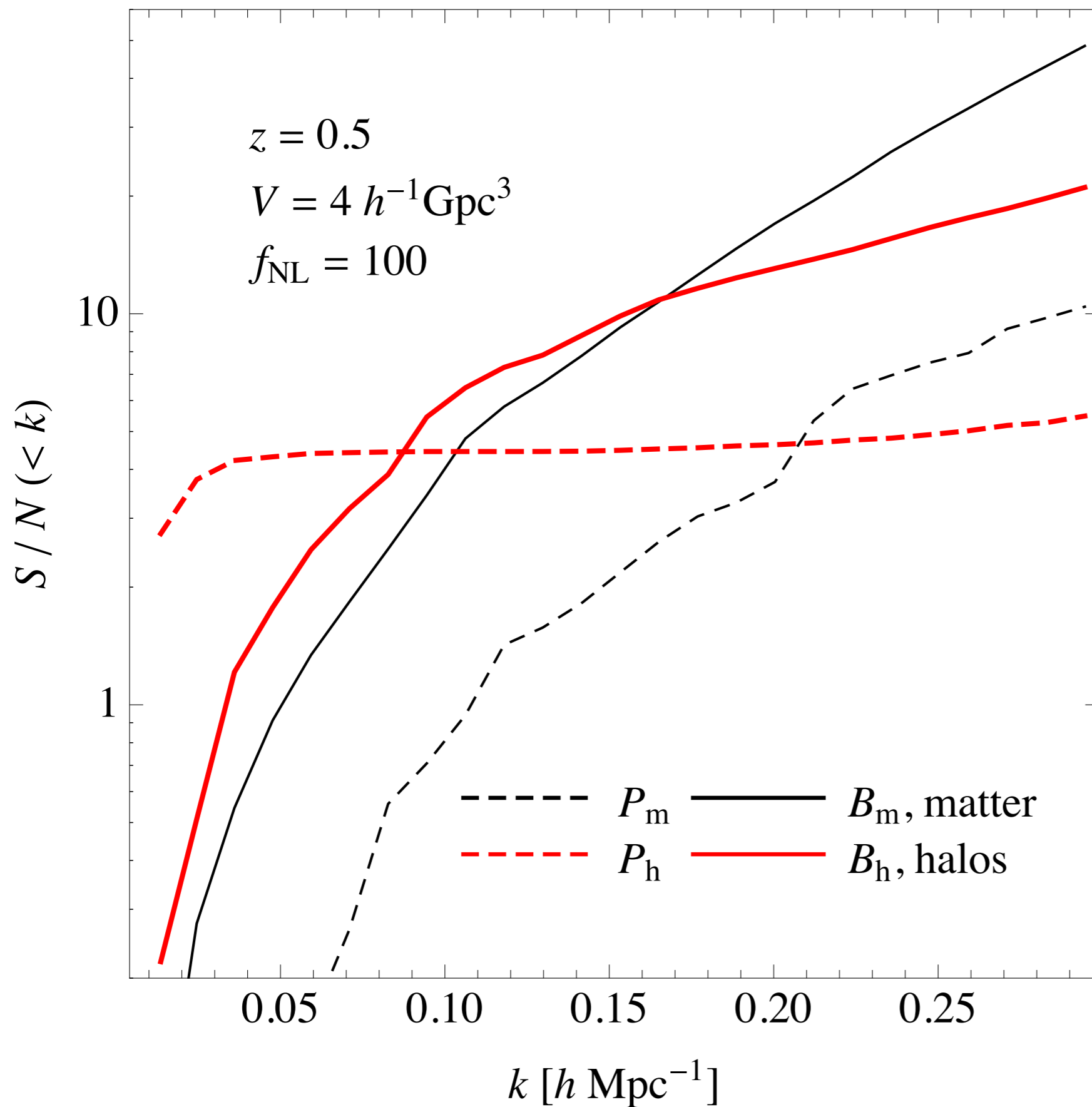
$$B_h(k_1, k_2, k_3) = b_1^3(f_{NL}, k) B(k_1, k_2, k_3) \\ + b_1(f_{NL}, k_1) b_1(f_{NL}, k_2) b_2(f_{NL}, k_1, k_2) P(k_1) P(k_2) + cyc.$$

Squeezed configurations $B(\Delta k, k, k)$
as a function of k with $\Delta k = 0.01 h/\text{Mpc}$



ES, Crocce & Desjacques (*in preparation*)

Power Spectrum vs. Bispectrum



Cumulative signal-to-noise for the effect of NG initial conditions on matter and galaxy correlators (P & B)

Sum of all configurations up to k_{max}

$$\left(\frac{S}{N}\right)_P^2 = \sum_k^{k_{\text{max}}} \frac{(P_{\text{NG}} - P_G)^2}{\Delta P^2}$$

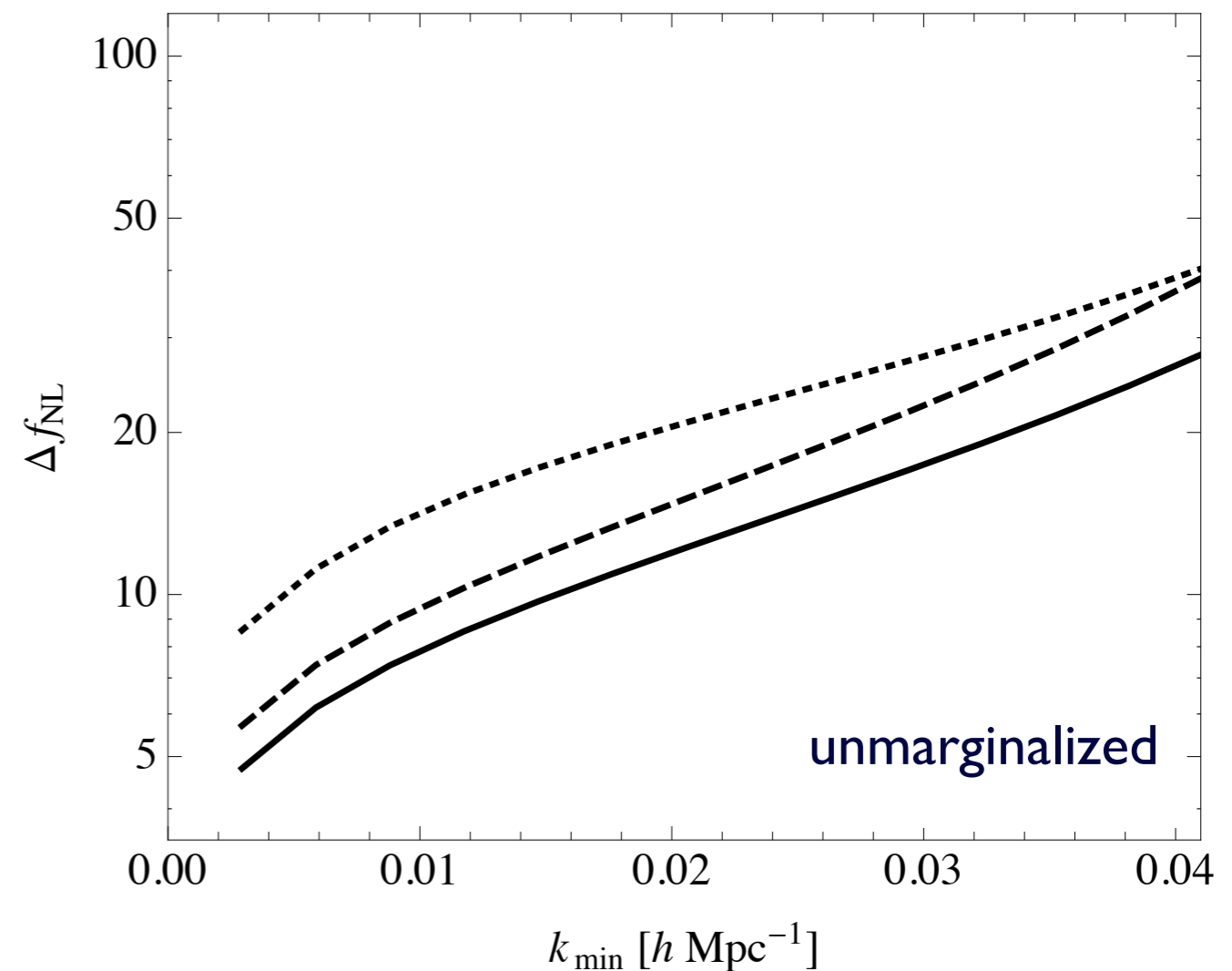
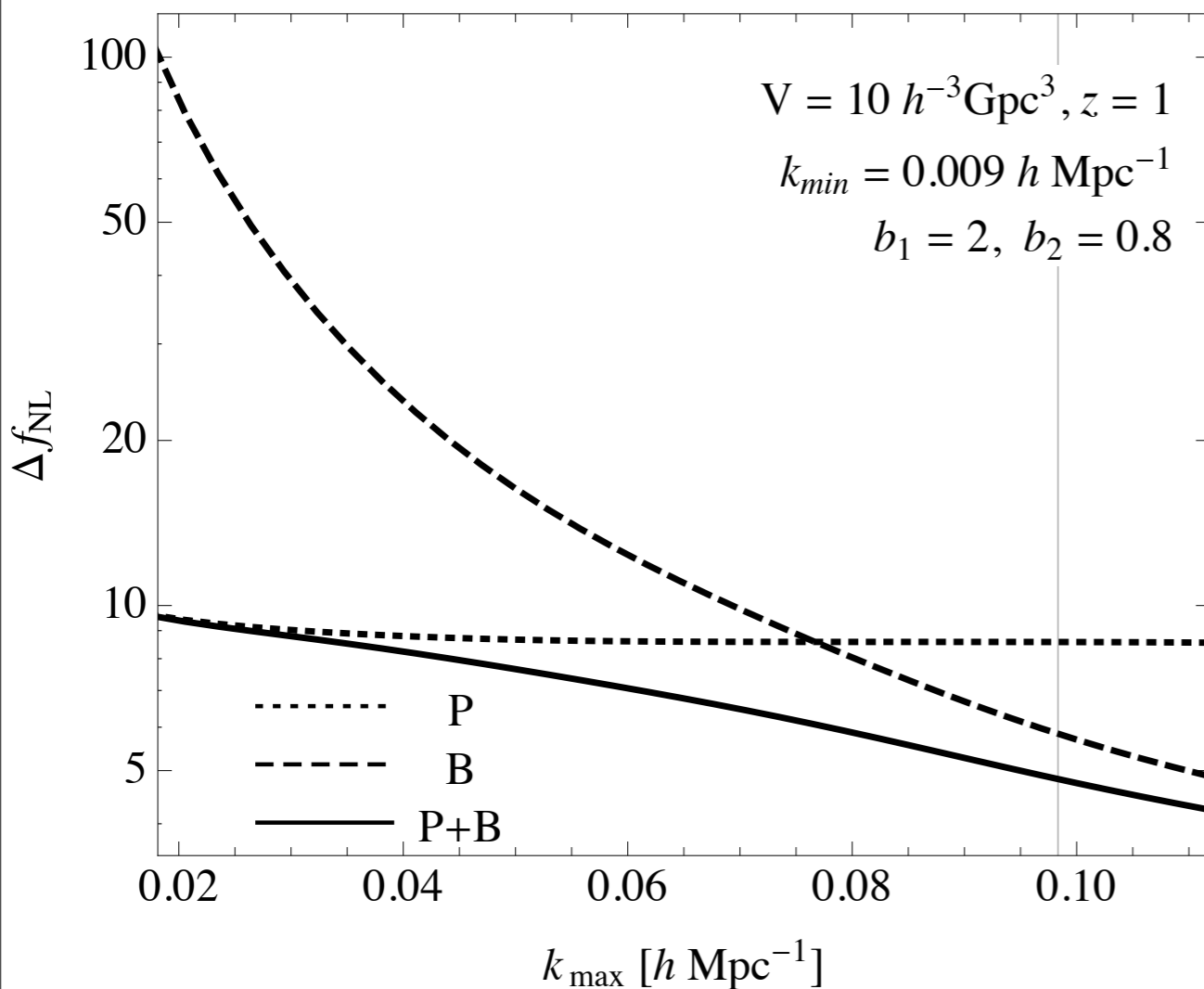
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An *unrealistic* Fisher matrix analysis

Assuming *perfect knowledge* of a complete galaxy population in a 10 Gpc^3 volume at redshift z with density 10^{-3} Mpc^{-3}

Ignoring *any* complication no matter how relevant and pertinent (covariance, redshift distortions, selection function, degeneracies, etc ...)

We can estimate the uncertainty on f_{NL} (local) from Power Spectrum & Bispectrum (& both)

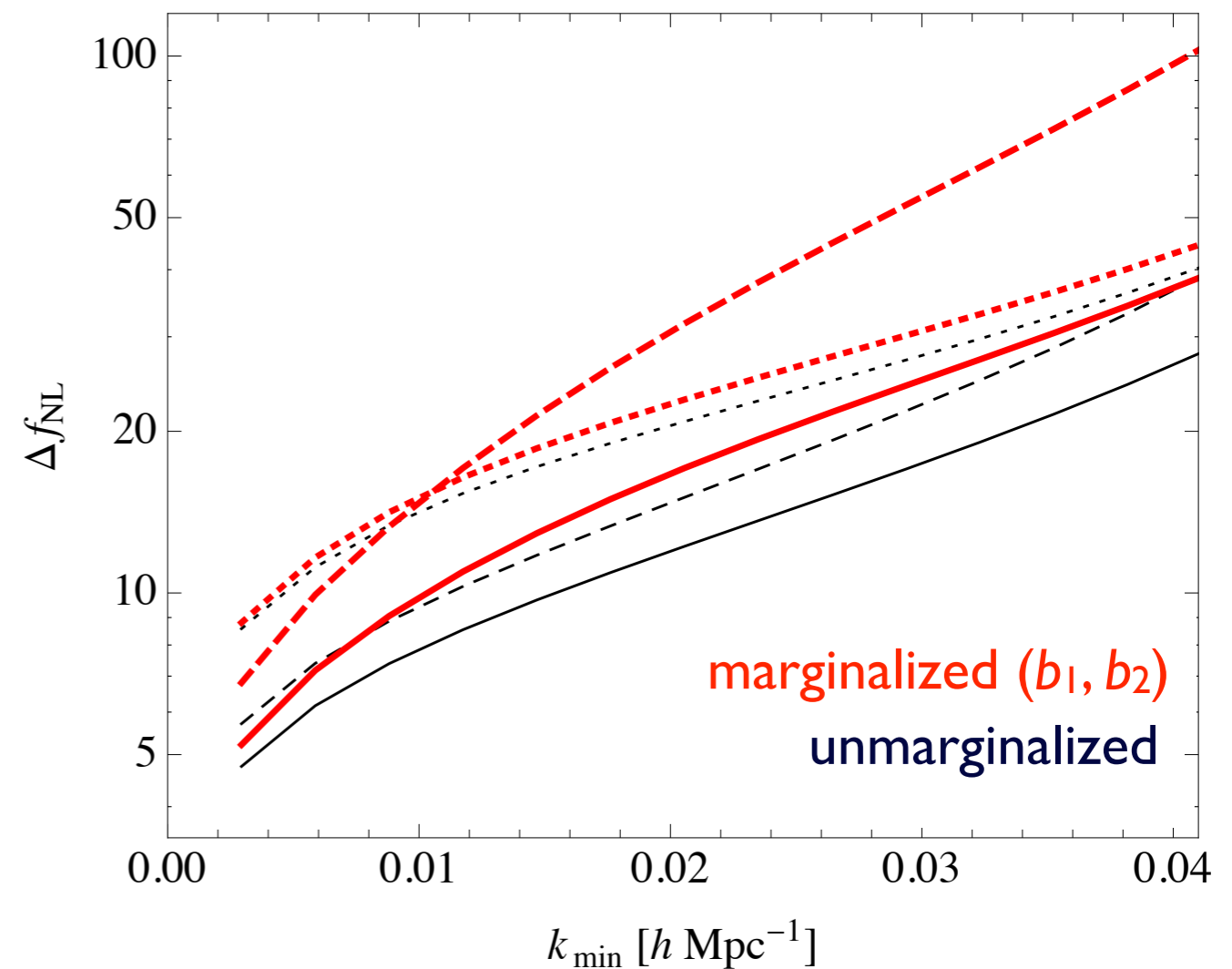
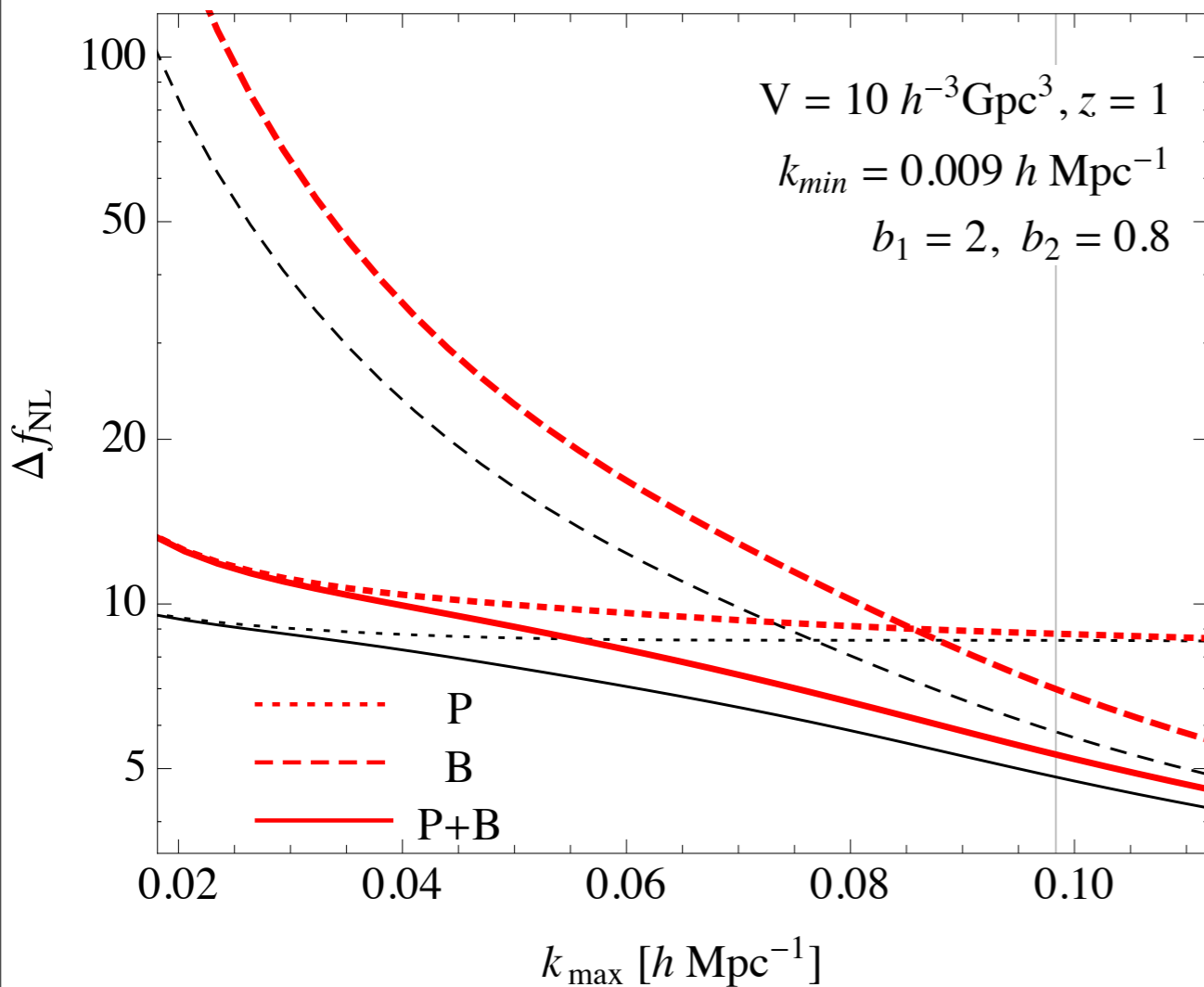


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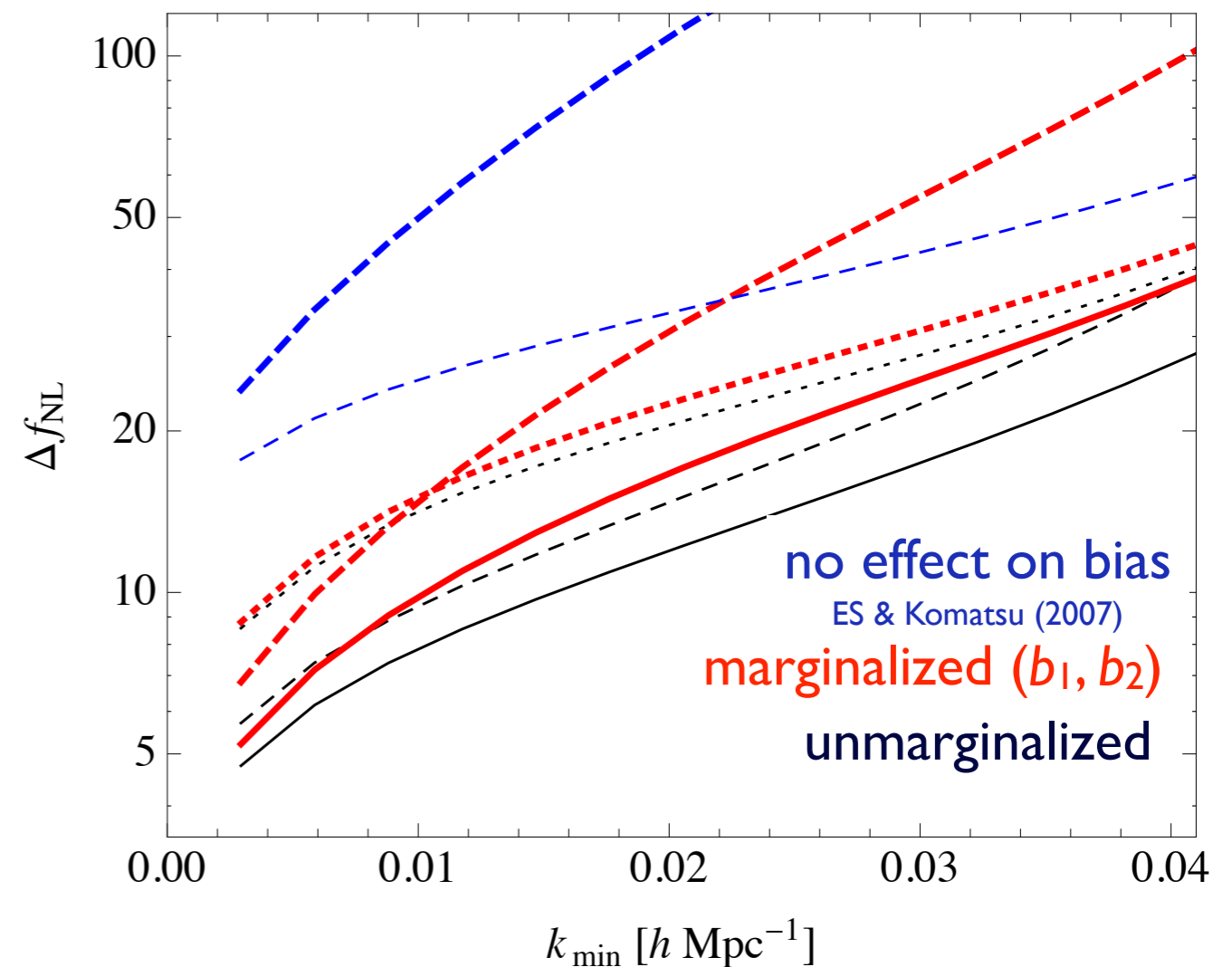
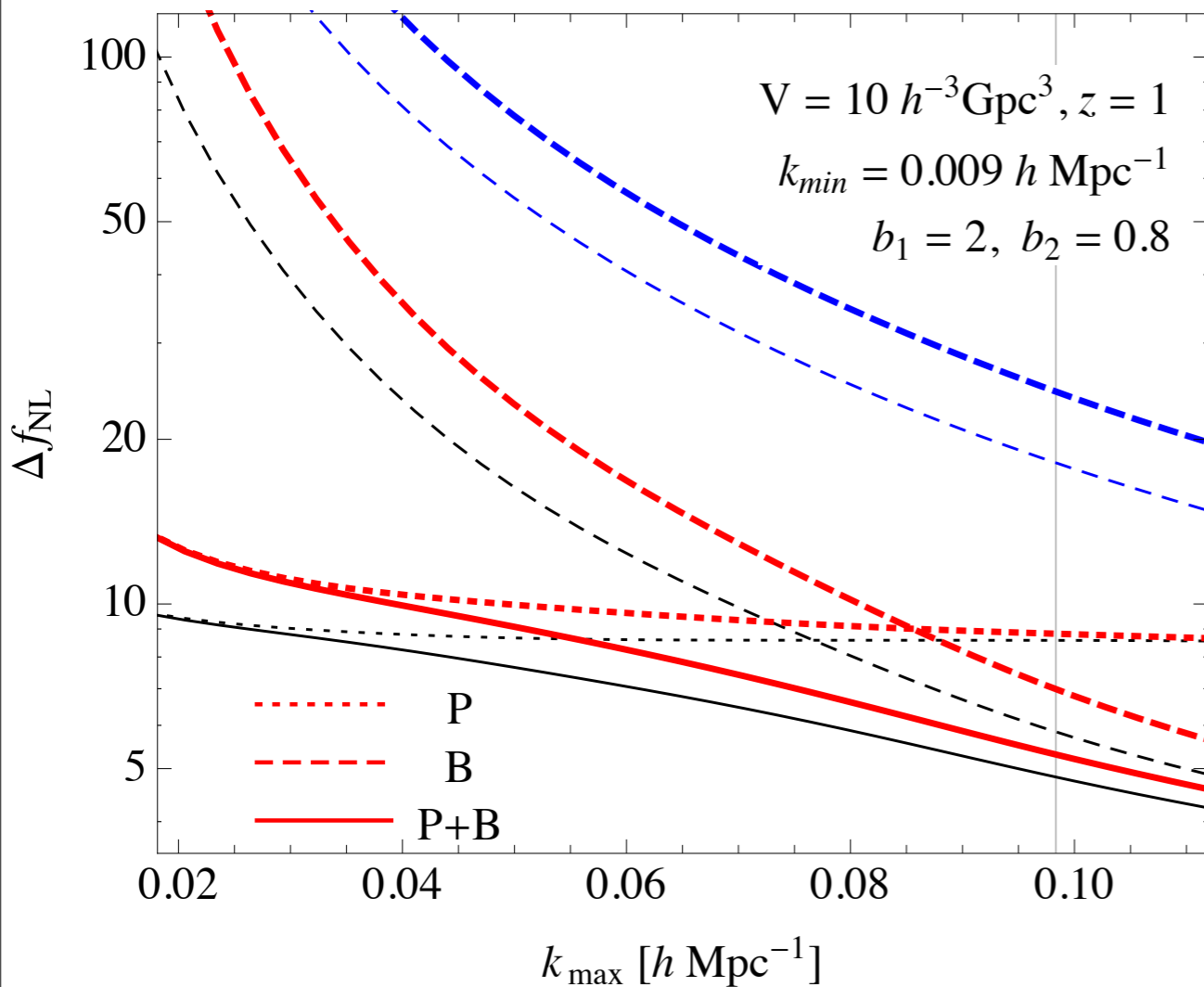


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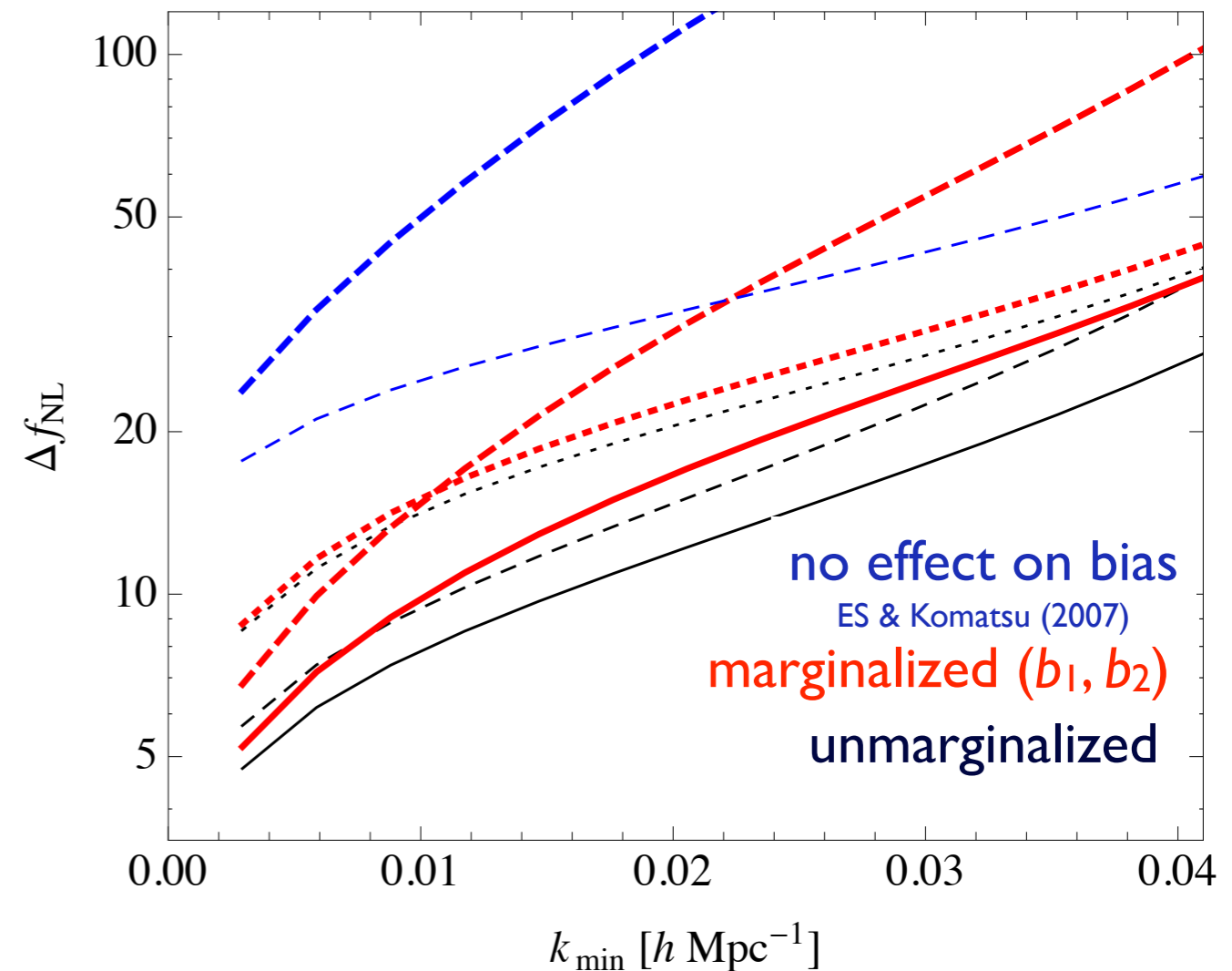
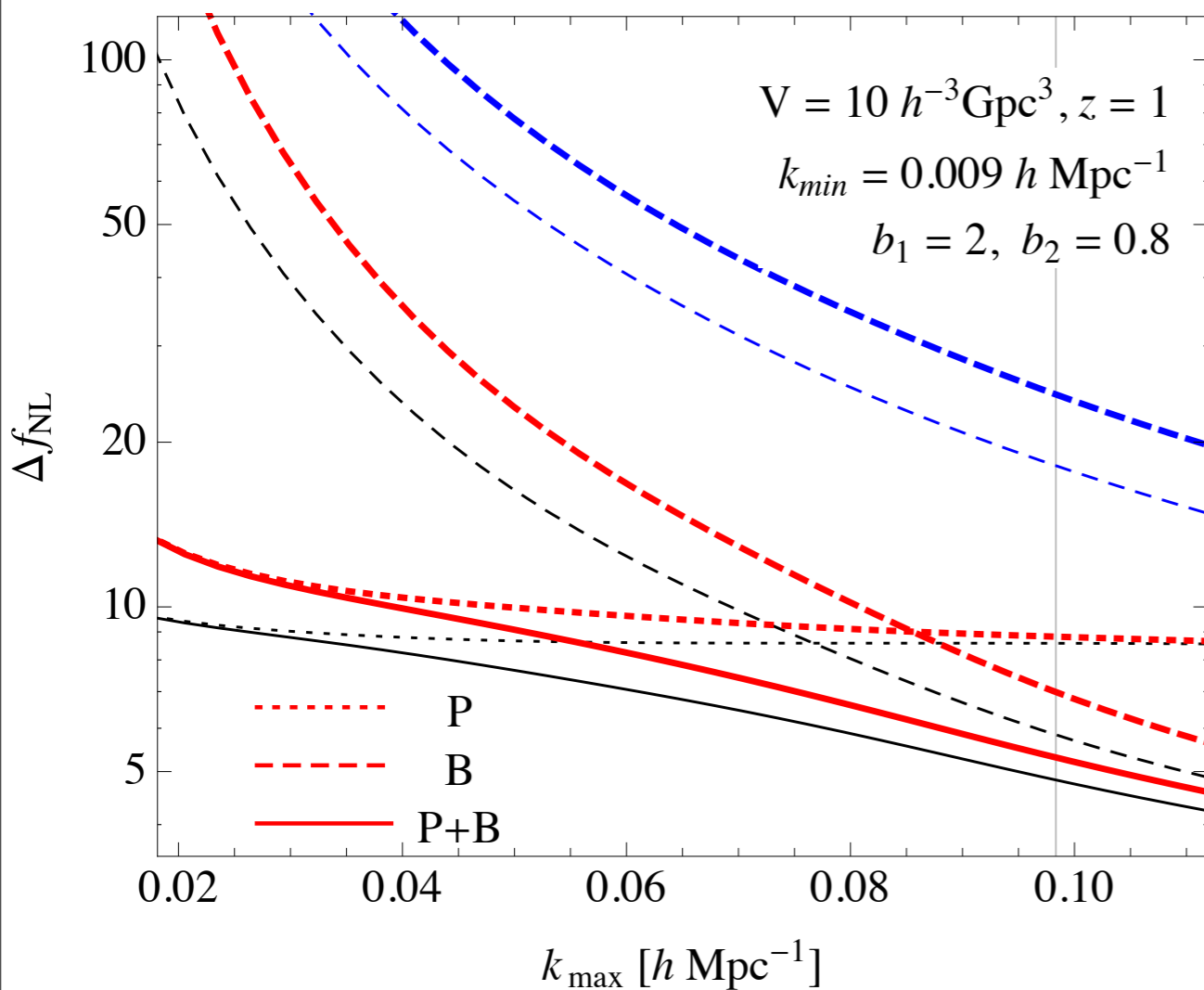
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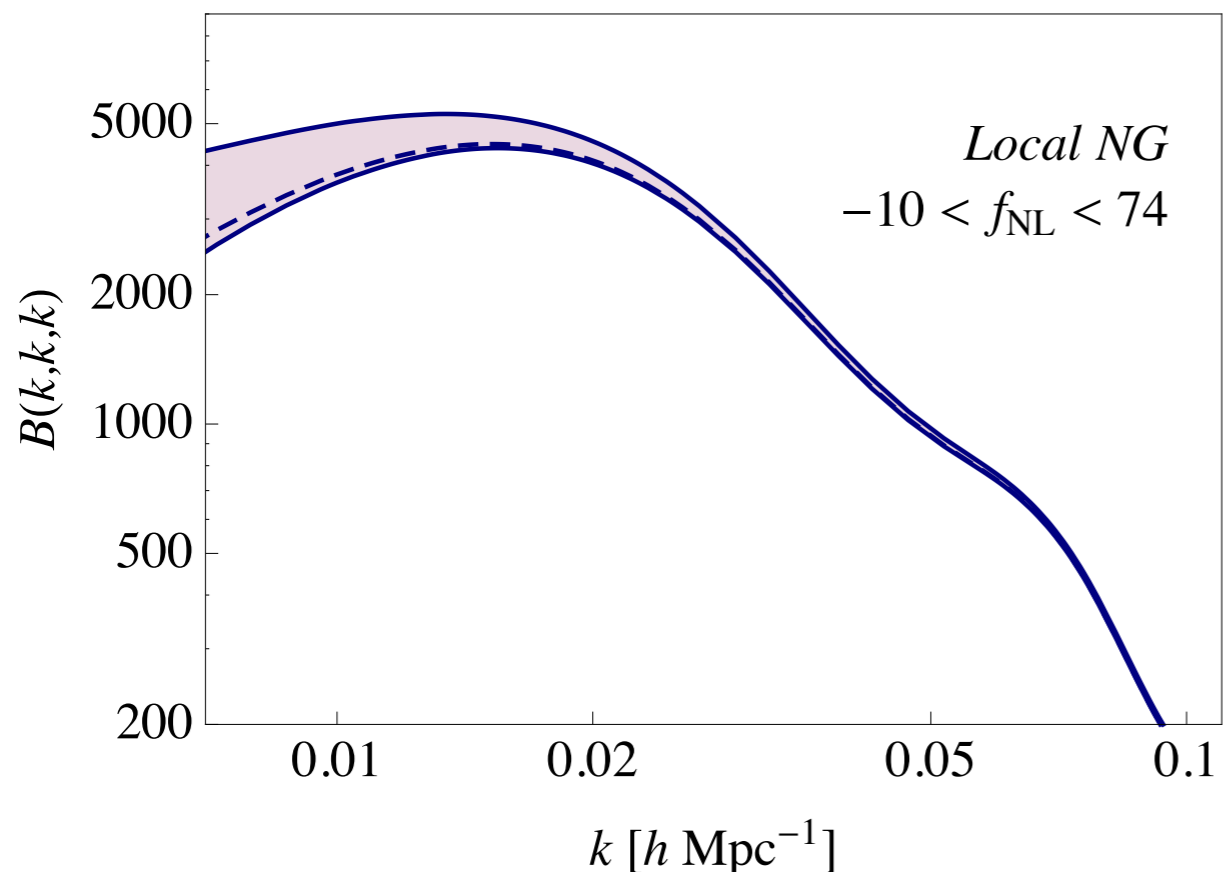


Conclusions

- We have a (relatively simple) model for the large scales galaxy bispectrum with local NG initial conditions
- Bispectrum measurements in LSS surveys can confirm and improve constraints on f_{NL} from the power spectrum (particularly for non-local models ...)

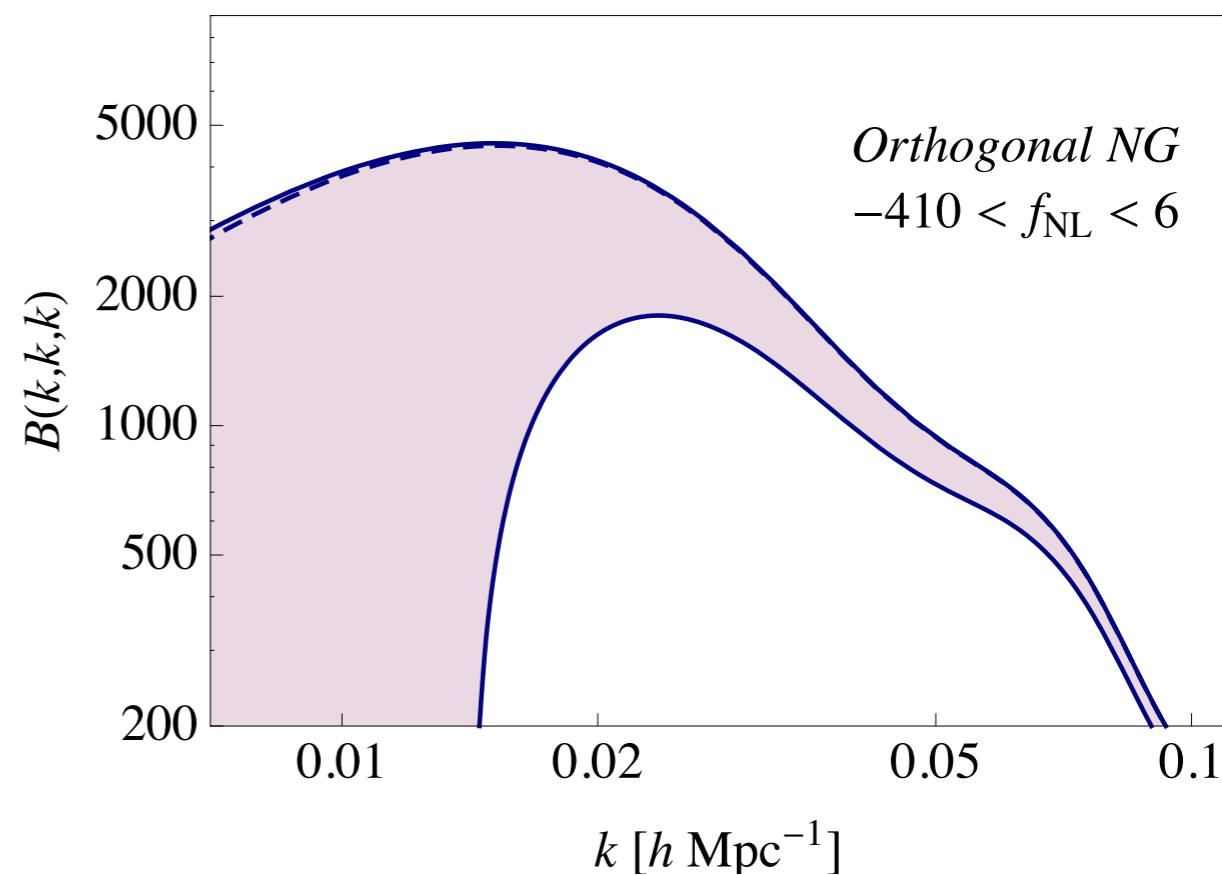
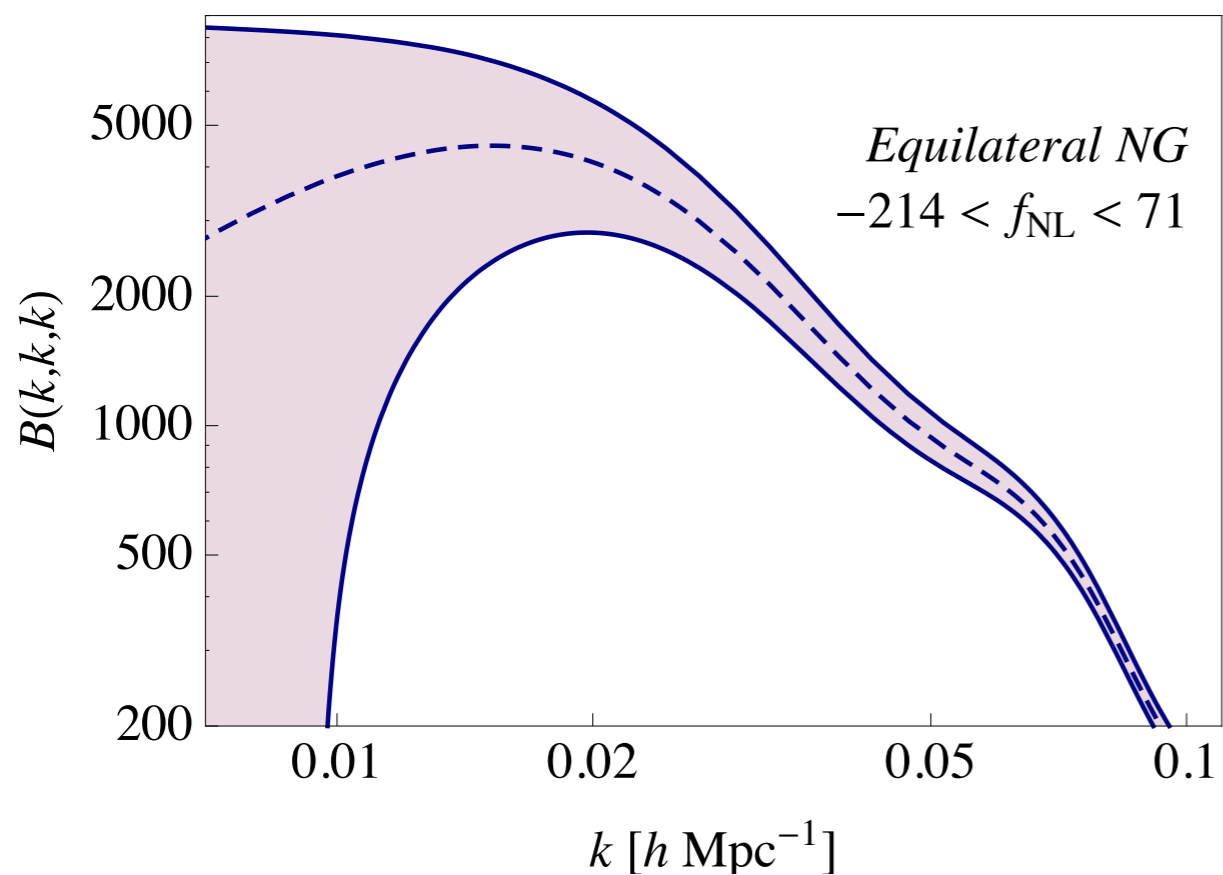


The matter bispectrum and PNG: *large scales*



Current CMB constraints for different models of non-Gaussianity as uncertainties on the equilateral configurations of the matter bispectrum

$$B \simeq B_0 + B_G^{\text{tree}} [P_0]$$

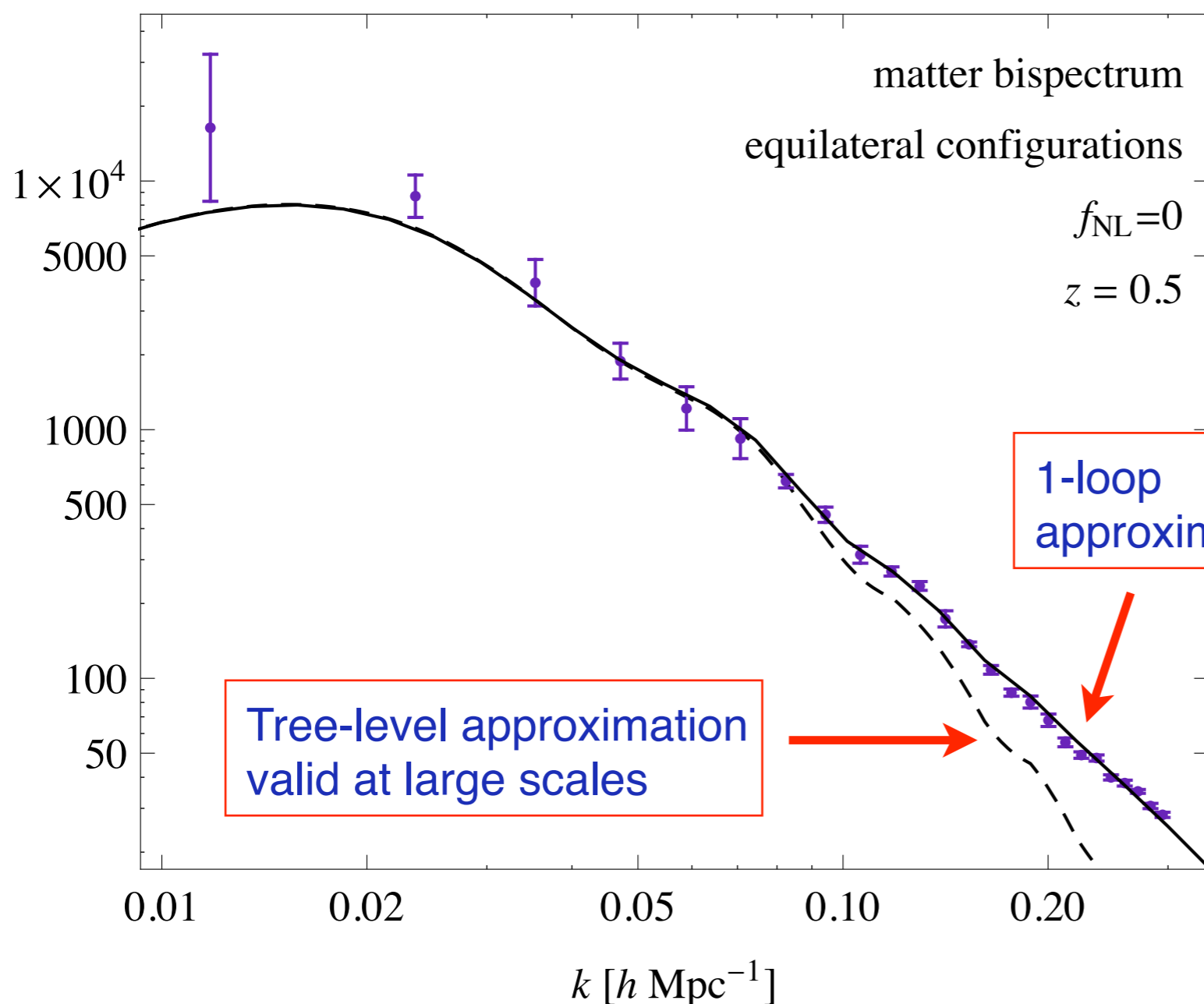


The Matter Bispectrum induced by Gravity

$$B_G = B_G^{tree}[P_0] + B_G^{loop}[P_0]$$

→ $B_G^{tree}(k_1, k_2, k_3) = 2 F_2(\vec{k}_1, \vec{k}_2) P_0(k_1) P_0(k_2) + 2 \text{ perm.}$

The bispectrum induced by gravity has a well defined dependence on scale and on the shape



The equilateral configurations of the matter bispectrum:

$B(k, k, k)$ vs. k

Numerical simulations and PT predictions

E.S., M. Crocce, & V. Desjacques (2010)

Non-Gaussianity from Gravitational Instability

At large scales fluctuations are small, $\sigma_\delta \ll 1$, even at low redshift we can study their evolution in terms of **Perturbation Theory**

Equations of motion for matter density and velocity:

$$\delta, \vec{v}$$

- Continuity eq. $\frac{\partial \delta}{\partial \tau} + \nabla \cdot [(1 + \delta)\vec{v}] = 0$
- Euler eq. $\frac{\partial \vec{v}}{\partial \tau} + \mathcal{H}\vec{v} + (\vec{v} \cdot \nabla)\vec{v} = -\nabla\phi$
- Poisson eq. $\nabla^2 \phi = \frac{3}{2}\mathcal{H}^2\Omega_m\delta$

Perturbative solution for the matter density, in Fourier space

$$\delta_{\vec{k}} = \delta_{\vec{k}}^{(1)} + \delta_{\vec{k}}^{(2)} + \dots$$

Linear solution

$$\delta_{\vec{k}}^{(2)} = \int d^3q F_2(\vec{k} - \vec{q}, \vec{q}) \delta_{\vec{k}-\vec{q}}^{(1)} \delta_{\vec{q}}^{(1)}$$

Quadratic nonlinear correction

Initial conditions

B_0 and T_0 vanish for Gaussian initial conditions!

$$\langle \delta_{\vec{k}_1}^{(1)} \delta_{\vec{k}_2}^{(1)} \rangle = \delta_D(\vec{k}_1 + \vec{k}_2) P_0(k_1)$$

$$\langle \delta_{\vec{k}_1}^{(1)} \delta_{\vec{k}_2}^{(1)} \delta_{\vec{k}_3}^{(1)} \rangle = \delta_D(\vec{k}_1 + \vec{k}_2 + \vec{k}_3) B_0(k_1, k_2, k_3)$$

$$\langle \delta_{\vec{k}_1}^{(1)} \delta_{\vec{k}_2}^{(1)} \delta_{\vec{k}_3}^{(1)} \delta_{\vec{k}_4}^{(1)} \rangle = \delta_D(\vec{k}_1 + \dots + \vec{k}_4) T_0(\vec{k}_1, \vec{k}_2, \vec{k}_3, \vec{k}_4)$$

Perturbative solution for the matter 3-point function

$$\langle \delta\delta\delta \rangle = \langle \delta^{(1)} \delta^{(1)} \delta^{(1)} \rangle + \langle \delta^{(1)} \delta^{(1)} \delta^{(2)} \rangle + \dots \quad \text{loop corrections}$$

= $B_0 = 0$ for Gaussian initial conditions

non-zero bispectrum induced by gravity!

Non-Gaussianity from **Galaxy Bias** (more problems?)

Additional non-Gaussianity in the **galaxy distribution** is induced by **nonlinear galaxy bias**

The relation between the **observed galaxy overdensity** and the matter density is nonlinear

$$\delta_g(x) \equiv \frac{n_g(x) - \bar{n}_g}{\bar{n}_g} = f[\delta(x)] \quad \text{local bias}$$

At large scales, we expand it in a Taylor series

$$\delta_g(x) = b_1 \delta(x) + \frac{1}{2} b_2 \delta^2(x) + \dots$$

Linear bias

Quadratic bias correction

Perturbative solution for the **galaxy 3-point function**

$$\langle \delta_g \delta_g \delta_g \rangle = b_1^3 \langle \delta \delta \delta \rangle + b_1^2 b_2 \langle \delta \delta \delta^2 \rangle + \dots$$

matter bispectrum

bispectrum induced by nonlinear bias

$$B_g(k_1, k_2, k_3) = b_1^3 B(k_1, k_2, k_3) + b_1^2 b_2 P(k_1) P(k_2) + 2 \text{ perm.} + \dots$$

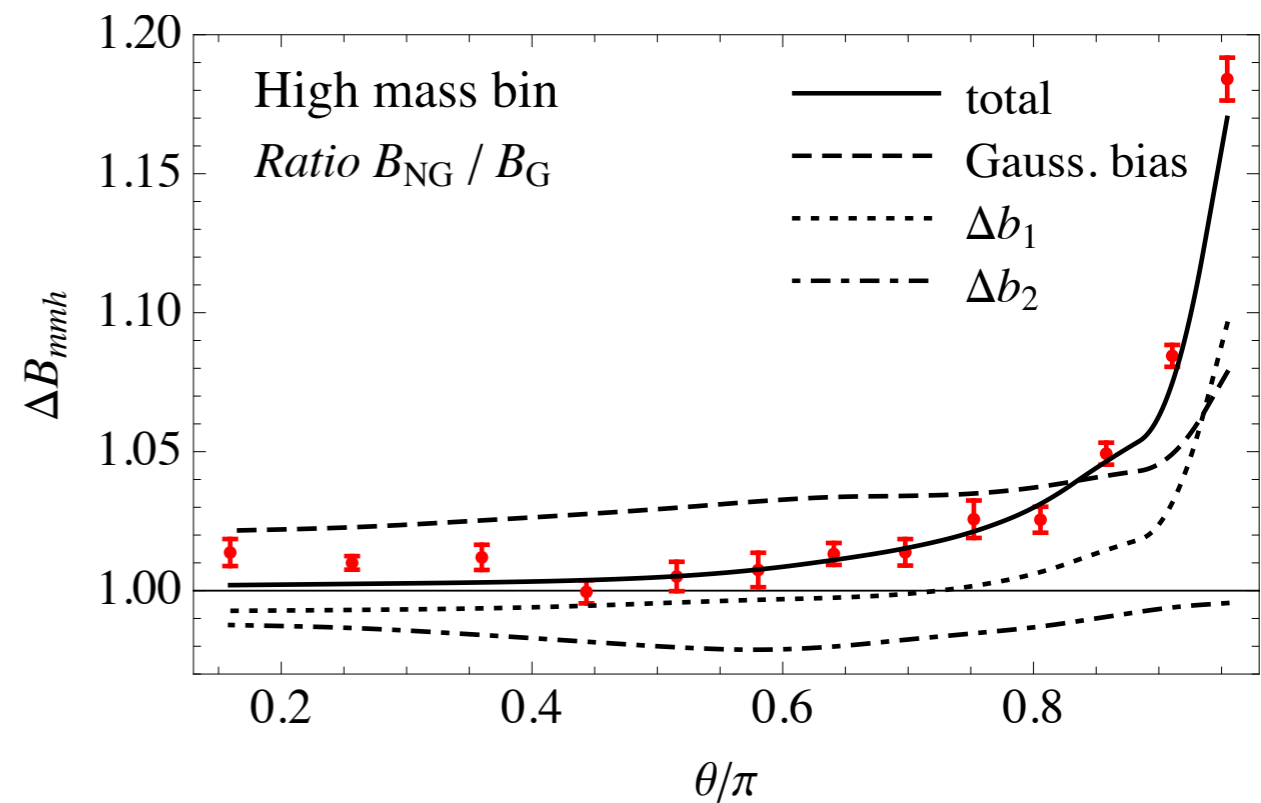
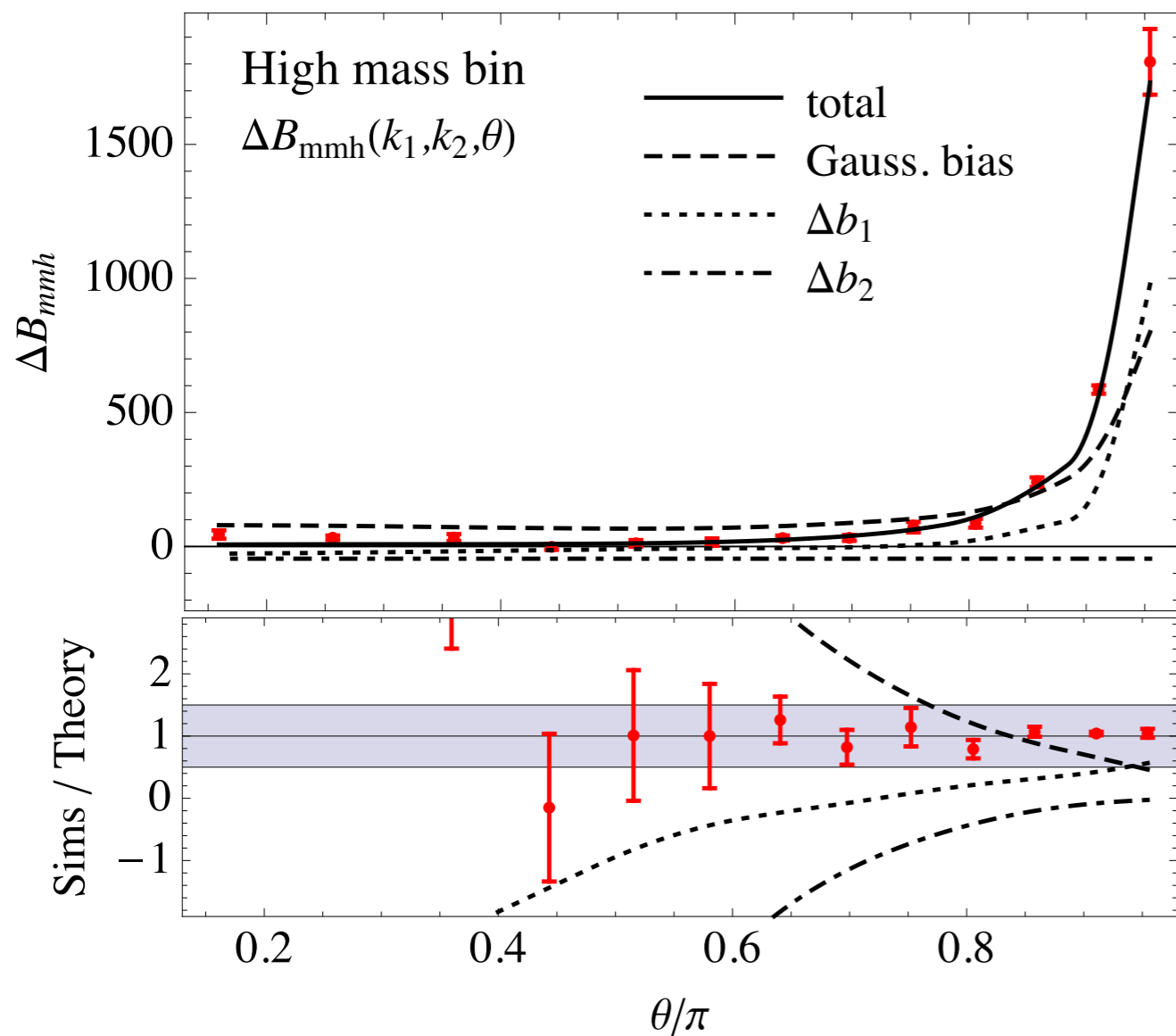
The component induced by bias has a different dependence on the shape of the triangle

Effects of PNG on the galaxy bispectrum

Clearly, the effect on galaxy bias affects as well the **galaxy bispectrum**

$$B_g(k_1, k_2, k_3) = \overset{b_{1,G} + \Delta b_{1,NG}(f_{NL}, k)}{\uparrow} b_1^3 B(k_1, k_2, k_3) + \overset{b_{2,G} + \Delta b_{2,NG}(f_{NL}, \vec{k}_1, \vec{k}_2)}{\uparrow} b_1^2 b_2 P(k_1) P(k_2) + 2 \text{ perm.} + \dots$$

Scale-dependent bias corrections



ES, Crocce & Desjacques (in preparation)

The matter bispectrum and PNG: *small scales*

